

On a submanifold of a manifold with areal metric (**)

Summary. Manifolds with areal metric have been defined and studied by various authors namely Davies [3], Kawaguchi [6], Martin [7] and Rund [9], [10] etc. Also the submanifolds of different type of manifolds have been studied by Bompiani [1], Davies [2], Douglas [4], Eliopoulos [5], Rund [8] and others, but no attempt has been made to study the submanifolds of the manifolds with areal metric, defined by Rund [10]. The purpose of the present paper is to study the submanifolds of the manifolds cited above. In this paper I have studied properties of the normal vectors, induced connection parameters, induced covariant derivative, normal curvatures of the submanifold etc. Finally Gauss-Codazzi equations in such manifolds have been obtained.

1. INTRODUCTION. — Let X_n be a differentiable manifold referred to local coordinates x^i ($i = 1, \dots, n$) such that

$$(1.1) \quad x^i = x^i(t^a), \quad (a = 1, \dots, l),$$

where t^a denotes a system of independent parameters of l -dimensional subspace C_l of X_n ($l < n$), such that Rund [10]:

$$(1.2) \quad \dot{x}_a^i = \partial x^i / \partial t^a.$$

Let us suppose that we are given a function $L(x^i, \dot{x}_a^i)$ of the $n + n l$ variables x^i, \dot{x}_a^i which satisfies:

I) the Lagrangian L is of class C^1 in all its arguments and it is a scalar with respect to the transforms of the local coordinates x^i of X_n ,

II) L is positive for all independent set of arguments \dot{x}_a^i ,

III) the integral

$$\int_{R_l} L(x^i, \dot{x}_a^i) dt^1 \dots dt^l,$$

where R_l is a finite simply connected region, is independent of the choice of parameters t^a of the subspace,

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IV the $n1 \times n1$ determinant

$$D = \det \left[\frac{1}{2} \frac{\partial^2 L^{n1}}{\partial \dot{x}_a^i \partial \dot{x}_b^j} \right],$$

is non-vanishing for linearly independent \dot{x}_a^i .

The third condition of the above can equivalently be expressed as

$$(1.3) \quad (\partial L_i / \partial \dot{x}_a^j) \dot{x}_b^j = \delta_{ab}^i L.$$

For such a manifold the metric tensor is defined by

$$(1.4) \quad g_{ij}^{ab}(\dot{x}^a, \dot{x}^b) = \frac{1}{2} \frac{\partial^2 L(\dot{x}^a, \dot{x}^b)}{\partial \dot{x}_a^i \partial \dot{x}_b^j}$$

which satisfies:

$$(1.5) \text{ a} \quad C_{ijk}^{ab} = \frac{1}{2} \left(\frac{\partial g_{ij}^{ab}}{\partial \dot{x}_k^c} \right)$$

and

$$(1.5) \text{ b} \quad g_{ab}^{ij} g_{ik}^{ab} = \delta_{jk}^i,$$

where

$$C_{ijk}^{ab} \dot{x}_a^i = 0 \quad \text{and} \quad C_{ijk}^{ab} \dot{x}_b^j \dot{x}_a^i = 0.$$

The coefficients of connection for the manifold are defined by Rund [10]:

$$(1.6) \quad \Gamma_{ki}^a = \Gamma^{-1} \left[P_{kij}^a + \frac{\partial P_{kij}^a}{\partial \dot{x}_b^c} \dot{x}_a^b \right],$$

where

$$P_{kij}^a = G_{ij}^{ab} (V_{ij,k}^{(a)} - C_{ijk}^{ab} \dot{x}_a^k),$$

$$V_{ij,k}^{(a)} = \frac{1}{2} \left(\frac{\partial g_{ij}^{ab}}{\partial \dot{x}^k} + \frac{\partial g_{ik}^{ab}}{\partial \dot{x}^j} - \frac{\partial g_{jk}^{ab}}{\partial \dot{x}^i} \right),$$

$$C_{kij}^{ab} = \frac{\partial C_{kij}^{ab}}{\partial \dot{x}_a^c}, \quad G_{ij}^{ab} = \frac{1}{2} (g_{ij}^{ab} + g_{ji}^{ab}).$$

Also these connection coefficients have the following properties:

$$(1.7) \text{ a} \quad \Gamma_{ki}^a \dot{x}_b^k = P_{kij}^a \dot{x}_a^k$$

and

$$(1.7) \text{ b} \quad \Gamma_{ki}^a \dot{x}_b^k \dot{x}_c^i = P_{kij}^a \dot{x}_a^k \dot{x}_c^i.$$

If X_{ϵ}^i is a set of linearly independent differentiable vector fields tangent to C , in X_{ϵ} , then we have

$$(1.8) \quad X_{\epsilon|j}^i = \frac{\partial X_{\epsilon}^i}{\partial \dot{x}^j} - \frac{\partial X_{\epsilon}^i}{\partial \dot{x}_a^j} \Gamma_{ik}^a \dot{x}_a^k + \Gamma_{ij}^a X_{\epsilon}^a,$$

where $X_{\epsilon|j}^i$ denotes the covariant derivative of X_{ϵ}^i with respect to \dot{x}^j .

Also we have

$$(1.9) \quad X_{\epsilon}^1|_{kb} - X_{\epsilon}^1|_{bk} = X_{\epsilon}^1 K_{\epsilon}^a|_{kb} - \frac{\partial X_{\epsilon}^1}{\partial x^k} K_{\epsilon}^a|_{kb} x^b + X_{\epsilon}^1|_{k} T_{\epsilon}^a|_{b},$$

where

$$(1.10) \quad K_{\epsilon}^a|_{kb} = \left(\frac{\partial \Gamma_{kj}^a}{\partial x^b} - \frac{\partial \Gamma_{jk}^a}{\partial x^k} \Gamma_{\epsilon}^i|_{kb} x^j \right) + \Gamma_{\epsilon}^i|_{kp} \Gamma_{\epsilon}^p|_{bj} - \left(\frac{\partial \Gamma_{kj}^a}{\partial x^k} - \frac{\partial \Gamma_{jk}^a}{\partial x^k} \Gamma_{\epsilon}^i|_{pk} x^j \right) - \Gamma_{\epsilon}^i|_{kp} \Gamma_{\epsilon}^p|_{bj},$$

$$(1.11) \quad T_{\epsilon}^a|_{kb} = \Gamma_{\epsilon}^a|_{kb} - \Gamma_{\epsilon}^a|_{bk}$$

and

$$(1.12) \quad K_{\epsilon}^a|_{jkb} = K_{\epsilon}^a|_{jkb} + T_{\epsilon}^a|_{kj} - T_{\epsilon}^a|_{jk} + T_{\epsilon}^a|_{jk} + (T_{\epsilon}^a|_{kj} T_{\epsilon}^i|_{kb} - T_{\epsilon}^a|_{ki} T_{\epsilon}^i|_{kb}).$$

2. SUBMANIFOLD OF X_n . — An m -dimensional submanifold X_m of X_n ($1 < m < n$) may be represented parametrically by the equations

$$(2.1) \quad x^i = x^i(u^A), \quad (A = 1, \dots, m),$$

where we suppose that the variables u^A form a coordinate system of X_m .

Now we shall parametrize the submanifold X_m by the vectors of C_1 such that

$$(2.2) \quad \hat{u}^A = \partial x^i / \partial u^A t^A, \quad (x = 1, \dots, l).$$

Along any coordinate curve of parameter u^A in X_m , the vector whose n components are $B^i_A = \partial x^i / \partial u^A$, where the matrix $\| B^i_A \|$ is of rank m , is tangential to the curve and corresponding to m independent variables u^A , there are m such linearly independent vector fields tangential to X_m , in terms of which any vector tangential to X_m is linearly expressible.

In particular, if $d x^i$ is a small displacement tangential to X_m , it follows that

$$(2.3) \quad d x^i = B^i_A d u^A,$$

where $d u^A$ denotes the same displacement in terms of the coordinates of X_m . Thus if we denote the components of a vector X_{ϵ}^i tangent to X_m by X_{ϵ}^i in terms of u^A system, we have

$$(2.4) \quad X_{\epsilon}^i = B^i_A X_{\epsilon}^A.$$

By virtue of equation (2.4) we can also obtain

$$(2.5) \quad \partial x^i / \partial u^A \hat{u}^A = B^i_A \hat{u}^A.$$

In case $x = \beta$ in (2.5) we shall express B^i_A by B^i_A . Now in analogy with the definition of the metric tensor for the manifold X_n we can also define the metric tensor for the manifold X_m as follows:

$$(2.6) \quad g_{\epsilon\delta}^a(u^c, \hat{u}^c_{\epsilon}) = \frac{1}{2} \left[\frac{\partial^2}{\partial \hat{u}^c_{\epsilon} \partial \hat{u}^c_{\delta}} L(u^c, \hat{u}^c_{\epsilon}) \right],$$

where $\bar{L}(u^c, \bar{u}_c^e)$ satisfies similar properties as $L(x^a, \dot{x}_a^b)$. Also for \bar{L} we have

$$(2.7) \quad (\partial \bar{L} / \partial \bar{u}_c^e) \bar{u}_c^e = \partial_5^e \bar{L}.$$

Since the metric of the manifold X_n induces a similar metric on the manifold X_m we can write

$$(2.8) \quad \bar{L}(u^c, \bar{u}_c^e) = L(x^i(u^c), \bar{u}_c^e B_i^e).$$

If we differentiate (2.8) successively with respect to \bar{u}_a^d and \bar{u}_b^d and use (1.4) and (2.6) we can establish a relationship between the two metric tensors in the following form

$$(2.9) \quad g_{AB}^{ab}(u^c, \bar{u}_c^e) = g_{ij}^{ab}(x^i, \dot{x}_i^b) B_i^A B_j^B.$$

As we can define the inverse of the metric tensor $g_{ij}^{ab}(x^i, \dot{x}_i^b)$ by $g_{ab}^{ij}(x^i, \dot{x}_i^b)$ such that Rund [10]:

$$(2.10) \quad g_{ab}^{ij}(x^i, \dot{x}_i^b) = \frac{1}{2} \frac{\partial^2 H(x^i, \dot{x}_i^b)}{\partial \dot{x}_i^a \partial \dot{x}_i^b},$$

where

$$H(x^i, \dot{x}_i^b) = L(x^i, \dot{x}_i^b),$$

we can also define inverse of the metric tensor of the submanifold X_m such that it satisfies:

$$(2.11) \quad g_{ab}^{AB}(u^c, z_c^e) = \frac{1}{2} \frac{\partial^2 \bar{H}(u^c, z_c^e)}{\partial z_c^A \partial z_c^B},$$

where

$$z_c^A = g_{AB}^{ab} \bar{u}_b^c = \partial \bar{L}^{AB} / \partial \bar{u}_c^A.$$

From these relations we can obviously have

$$(2.12) \text{ a) } B_i^A = g_{ab}^{AD} g_{Dc}^{ab} B_c^i$$

and

$$(2.12) \text{ b) } B_i^A B_c^i = \delta_c^A.$$

3. NORMAL VECTORS. — A covariant vector Y_i^T is said to be normal to X_m at a point P, if it satisfies

$$(3.1) \quad Y_i^T B_i^A = 0.$$

Since the rank of the matrix $[B_i^A]$ is assumed to be m , it follows that there exist $l(n-m)$ linearly independent vectors N_μ^T , ($\mu = m+1, \dots, n$) normal to X_m . These normals may be chosen in the multiply infinite number of ways:

$$(3.2) \quad N_\mu^T B_i^A = 0.$$

With respect to a given direction \hat{x}_α^i in the tangent space of X_n we may choose a set of normals satisfying the relations:

$$(3.3) \text{ a} \quad N_\mu^i = g_{\alpha\beta}^i(x^h, x_\beta^h) N_\mu^j \hat{x}_\alpha^j,$$

and

$$(3.3) \text{ b} \quad g_{\alpha\beta}^i(x^h, \hat{x}_\alpha^h) N_\mu^i N_\nu^j = \delta_{\alpha\nu}^i.$$

As a consequence of (3.3) it can be noted that

$$(3.4) \quad B_\mu^h N_\mu^i(x^h, y_\beta^h) = 0.$$

Now we shall define a relation

$$(3.5) \quad \psi_j^i = B_\lambda^i B_j^\lambda - \delta_j^i,$$

which by virtue of equation (2.12) b leads to

$$\psi_j^i B_\nu^j = 0,$$

which implies that ψ_j^i is of the form $\sum_\mu N_\mu^i N_\mu^j \lambda_\mu^i$, where the factors λ_μ^i are given by

$$\psi_j^i N_\nu^i = \lambda_\nu^i, \quad (\nu = m+1, \dots, n).$$

But from equations (3.4) and (3.5) we have

$$(3.6) \quad \psi_j^i N_\nu^i = -N_\nu^j,$$

which implies

$$(3.7) \quad B_\lambda^i B_j^\lambda = \delta_j^i - N_j^i(x^h, \hat{x}_\beta^h),$$

where

$$(3.8) \quad N_j^i(x^h, \hat{x}_\beta^h) = \frac{d\alpha_i}{d\alpha_j} \sum_{\mu=m+1}^n N_\mu^i N_\mu^j(x^h, \hat{x}_\beta^h).$$

The immediate consequence of the above and the preceding formulae is the relation

$$(3.9) \quad g_{ij}^{\alpha\beta}(x^h, \hat{x}_\alpha^h) = g_{\alpha\beta}^i(u^D, \hat{u}_\alpha^D) B_i^A B_j^A + N_{ij}^{\alpha\beta}(x^h, \hat{x}_\alpha^h),$$

and

$$(3.10) \quad g_{\alpha\beta}^i = g_{\alpha\beta}^{AC} B_i^A B_C^i + N_{\alpha\beta}^i.$$

Furthermore we can also obtain

$$(3.11) \quad \frac{1}{2} \frac{\partial g_{AB}^i(u^D, \hat{u}_\alpha^D)}{\partial \hat{u}_\gamma^k} = C_{AB\gamma}^{\alpha\beta}(u^D, \hat{u}_\alpha^D) = C_{B\alpha\gamma}^{\alpha\beta}(x^h, \hat{x}_\alpha^h) B_A^i B_\beta^A B_\gamma^i,$$

which by virtue of (2.12) gives rise to

$$(3.12) \quad \frac{\partial B_i^A}{\partial u_\alpha^C} = 2 g_{57}^{AP} \{ C_{0b}^{57} B_c^b B_D^b - C_{kb}^{57} B_D^b B_c^k B_1^k \} .$$

Applying equation (3.7) to the last term of (3.12) we obtain after simplification

$$(3.13) \quad \frac{\partial B_i^A}{\partial u_\alpha^C} = 2 C_{0b}^{57} N_j^b B_k^A B_0^b$$

or alternatively

$$(3.13) \text{ a} \quad \frac{\partial B_i^A}{\partial u_\alpha^C} = 2 N_j^b B_0^b B_D^k g_{57}^{AD} C_{0k}^{57} .$$

From equation (3.7) and (3.13) a we deduce that

$$(3.14) \quad \frac{\partial N_j^b}{\partial u_\alpha^C} = - 2 g_{57}^{AD} B_A^b B_0^k B_D^k C_{kb}^{57} N_j^b ,$$

which together with (3.13) a implies

$$(3.15) \quad \frac{\partial B_i^A}{\partial u_\alpha^C} \dot{u}_\alpha^C = \frac{\partial N_j^b}{\partial u_\alpha^C} \dot{u}_\alpha^C = 0 .$$

So far we have discussed the general properties of normals, but now we shall see that in general we have two sets of normals to X_m at a point P of X_m . The first set of normals n_a^i , which are independent of the directions \dot{x}_0^b , satisfy

$$(3.16) \quad n_a^i B_A^i \equiv g_0^{5b} n_b^j B_A^j = 0 .$$

The solutions of the equations (3.16) are normalized by means of the relation

$$(3.17) \quad L(x, n) = 1 \text{ r } g_0^{5b} (x^b, n_0^b) n_a^i n_b^j = 1 .$$

The second set of normals can be defined by the solutions $n_a^{*i} (x^b, \dot{x}_0^b)$ of the equations

$$(3.18) \quad g_0^{5b} (x^b, \dot{x}_0^b) B_A^i n_a^{*i} (x^b, \dot{x}_0^b) = 0$$

and its solutions are normalized by

$$(3.19) \quad g_0^{5b} (x^b, n_0^{*b} (x^c, \dot{x}_0^c)) n_a^{*i} n_b^{*j} = 1 .$$

The normals n_a^{*i} will be called secondary normals. From equations (3.17) and (3.19) it follows that n_a^{*i} is proportional to $n_a^i (x^b, n_0^b)$ and therefore we can define tensors $\gamma_{0iAD}^{\alpha\beta} (u^C)$, ($\mu = m + 1, \dots, n$) which are independent of the directions as follows :

$$(3.20) \quad \gamma_{0iAD}^{\alpha\beta} = g_0^{5b} (x^b, n_0^b) B_A^i B_D^j .$$

We also define the following set of inverse parameters corresponding to B^i_A :

$$(3.21) \quad \gamma_{(a)1}^A(x^b) = g_0^{ab} (x^k, n^c) \gamma_{(a)25}^{AB} B^j_b,$$

so that we have

$$(3.22) \quad n^a_{(2)} B^i_a = 0, \quad \gamma_{(a)1}^A n^a_{(2)} = 0,$$

and

$$(3.23) \quad \gamma_{(a)1}^A B^i_c = \delta^i_c.$$

4. INDUCED CONNECTION PARAMETERS. — Let $x^i = x^i(s)$ be a curve C of X_m which is continuous and continuously differentiable, then if we change the parameter from t to s and the dot by dash we can say that $x^i_{(s)}$ is tangential to X_m . Let us consider a continuous and continuously differentiable vector field tangent to X_m ; which satisfies (2.4), then the induced covariant derivative of the vector field along C in the space X_m is defined by

$$(4.1) \quad X^i_{(s)1} = \frac{\delta X^i_A}{\delta u^A} - \frac{\delta X^i_E}{\delta u^E} \Gamma^E_{DB} \dot{x}^D + \Gamma^i_{EC} X^C_E,$$

where Γ^i_{EC} are induced connection parameters.

The tensor $X^i_{(s)1}$ defined by (4.1) is given by the projection onto X_m of the covariant derivative $X^i_{(s)1}$ of X^i_E with respect to X_m . Hence

$$(4.2) \quad g_0^{ab} B^i_A B^j_k X^i_{(s)1} = g_{AB}^0 X^D_{(s)1}.$$

If T^i_{AB} is some mixed tensor, its intrinsic derivative will be given by

$$(4.3) \quad \frac{D T^i_{AB}}{D s^3} = \left(\frac{\delta T^i_{AB}}{\delta u^C} - \frac{\delta T^i_{AB}}{\delta u^D} \Gamma^D_{EC} \dot{u}^E + \Gamma^i_{AC} T^k_{AB} - \Gamma^i_{AC} T^i_{Bk} \right) \frac{d u^C}{d s^3},$$

where Γ^i_{EC} is mixed connection.

Now applying equation (4.3) to the well known relation $x^b_{(s)} = B^b_A u^A_{(s)}$ we obtain on simplification

$$(4.4) \quad \frac{D x^i_{(s)}}{D s^3} = B^i_{AC} \dot{u}^A \dot{u}^C + B^i_A \frac{D \dot{u}^A}{D s^3},$$

where

$$(4.5) \quad B^i_{AC} = \frac{\delta^3 x^i}{\delta u^A \delta u^C} - \frac{\delta^3 x^i_E}{\delta \dot{u}^A \delta \dot{u}^E} \Gamma^E_{DC} \dot{u}^D + \Gamma^i_{AC} B^k_A B^k_C - B^i_D \Gamma^D_{AC}.$$

The expression B^i_{AC} which is a tensor may be considered as the generalised covariant derivative of B^i_A with respect to u^C . Because of the second term on the right hand side it can be seen that the tensor B^i_{AC} is not symmetric in its lower indices.

Now multiplying equation (4.2) by u_γ^c , we obtain

$$g_{ij}^{a\beta} B_{\alpha}^i \frac{D x_{\beta}^j}{D s^{\theta}} = \kappa_{AC} \frac{D \hat{u}_{\beta}^j}{D s^{\theta}},$$

which is satisfied by the tangent vector u_{γ}^c to any curve C in X_m .

Now in case of a submanifold X_m we define a curve C to be a geodesic G , if it satisfies

$$D \hat{u}_{\beta}^j / D s^{\theta} = 0,$$

hence obviously from the last relation we can get

$$g_{ij}^{a\beta} (x^k, \hat{x}_{\gamma}^b) B_{\alpha}^i \left(\frac{D x_{\beta}^j}{D s^{\theta}} \right)_G = 0.$$

Since the vector $D x_{\beta}^j / D s^{\theta}$, which defines a principal normal to a geodesic G , satisfies equation (3.18), it belongs to the space spanned by n_{α}^{*j} , therefore we have:

Theorem (4.1). — *The principal normal of a geodesic G of X_m lies in the space spanned by the secondary normals n_{α}^{*j} .*

Since we know that

$$g_{ij}^{a\beta} (x^k, \hat{x}_{\beta}^b) B_{\alpha}^i B_{\gamma}^j \Gamma_{DE}^C = g_{ij}^{a\beta} B_{\alpha}^i \left(\frac{\partial^2 x^j}{\partial u^D \partial u^E} - \frac{\partial^2 x_{\beta}^j}{\partial u_{\beta}^D \partial u_{\beta}^E} \Gamma_{BE}^C \hat{u}_{\gamma}^{\beta} + \Gamma_{BE}^C B_D^{\beta} B_{\beta}^k \right),$$

therefore by virtue of (4.5) we obtain

$$(4.7) \quad g_{ij}^{a\beta} (x^k, \hat{x}_{\beta}^b) B_{\alpha}^i B_{\gamma}^j = 0,$$

which implies:

Theorem (4.2). — *The tensor B_{AC} lies in the space spanned by the secondary normals n_{α}^{*j} .*

5. NORMAL CURVATURES. — Since B_{AC} lies in the space spanned by secondary normals it can be expressed as a linear combination of n_{α}^{*j} , thus

$$(5.1) \quad B_{AC}^i = \sum_{(\alpha)} \Omega_{\alpha}^{*i} (u^D, \hat{u}_{\beta}^D) n_{\alpha}^{*j},$$

where Ω_{α}^{*i} is called secondary second fundamental tensor of the submanifold X_m . Also this tensor is not symmetric in A and C .

Multiplying equation (5.1) by n_{α}^{*j} and putting

$$(5.2) \quad \sum_{\mu} \Omega_{\mu}^{*i} \cos (n, n^{\mu}) = \Omega_{\alpha}^{*i},$$

we find

$$(5.3) \quad D_{(C)}^j n_j^b = \Omega_{(C)AC}^b .$$

The tensor $\Omega_{(C)AC}^b$ is also called the second fundamental tensor and in (5.2) $\cos(n, n^*)$ is the cosine of the angle between the two types of normals and is expressed as

$$(5.4) \quad \cos(n, n^*) = \frac{g_{ij}^{(C)}(x, n) n_i^a n_j^b}{[g_{ij}^{(C)}(x, n) n_i^a n_j^b \cdot g_{ab}^{(C)}(x, n^*) n_r^a n_s^b]^{1/2}} .$$

Multiplying equation (5.3) by $u_a^A u_V^C$ and using

$$n_i^b \frac{d x_i^j}{d s^j} = n_i^b \frac{\lambda^k x^j}{\lambda u^A \lambda u^C} \dot{u}_a^A \dot{u}_V^C ,$$

we get on simplification

$$(5.5) \quad \Omega_{(C)AC}^b \dot{u}_a^A \dot{u}_V^C = n_i^b \frac{D x_i^j}{D s^j} .$$

Equation (5.5) holds for all curves of X_m with tangent vector x_a^A , but depends on the choice of L . If we differentiate the relation

$$n_i^a x_a^A = 0 ,$$

we find

$$\frac{D n_i^a}{D s^j} x_a^A = - n_i^a \frac{D x_a^A}{D s^j} ,$$

which implies that $n_i^a D x_a^A / D s^j$ depends on the choice of line element.

Using (5.5) in the identity

$$n_i^a \frac{D x_a^A}{D s^j} = \left| \frac{D x_a^A}{D s^j} \right| \left| n_i^a \right| \cos \left(n_i^a, \frac{D x_a^A}{D s^j} \right) ,$$

we obtain

$$(5.6) \quad \left| \frac{D x_a^A}{D s^j} \right| = \frac{\Omega_{(C)AC}^b \dot{u}_a^A \dot{u}_V^C}{\cos \left(n_i^a, \frac{D x_a^A}{D s^j} \right) \left| n_i^a \right|} ,$$

which gives an expression for the curvature of a curve of X_m .

Since $\Omega_{(C)AC}^b \dot{u}_a^A \dot{u}_V^C$ is same for all curves of X_m , tangent to x_a^A , equation (5.6) gives a generalisation of the Meuniers theorem of classical differential geometry.

Therefore we may regard

$$(5.7) \quad \frac{1}{R_{(\alpha)}^*} \frac{d\Omega}{ds} = \Omega_{(\alpha)AC}^* \dot{u}_\alpha^A \dot{u}_\gamma^C,$$

as normal curvature corresponding to the normals $n_{(\alpha)}^*$.

So far we have shown that to each direction at a point P of X_m , correspond l ($n-m$) normal curvatures associated with the given direction u_α^A , it can be easily proved in analogy to Finsler spaces Rund [8], that the principal directions will be given by the extreme values of $g_{AC}^*(u^h, u_\alpha^h) u_\alpha^A u_\gamma^C$, where u^h is kept fixed. In other words we can say that principal directions are those for which normal curvatures assume extreme values.

Now we shall define a secondary normal curvature associated to a line element (x^h, x_α^h) and depending on Ω^* . Let us consider

$$(5.8) \quad \frac{D x_\alpha^i}{D s^2} = \sum_{\mu} \lambda_{(\alpha)\mu} n_\mu^i + R_A^i \frac{D \dot{u}_\alpha^A}{D s^2},$$

for an arbitrary curve of X_m . If we multiply (5.8) by $n_{(\alpha)}^*$, we get

$$(5.9) \quad n_{(\alpha)}^* \frac{D x_\alpha^i}{D s^2} = \sum_{\mu} \lambda_{(\alpha)\mu} \cos(n, n^*),$$

which by virtue of (5.6) implies

$$(5.10) \quad \Omega_{(\alpha)AC}^* \dot{u}_\alpha^A \dot{u}_\beta^C = \sum_{\mu} \lambda_{(\alpha)\mu} \cos(n, n^*).$$

Thus we have

$$\lambda_{(\alpha)\mu} = \Omega_{(\alpha)AC}^* \dot{u}_\alpha^A \dot{u}_\beta^C$$

and hence for a geodesic (5.8) yields

$$(5.11) \quad \frac{D x_\alpha^i}{D s^2} = \sum_{\mu} \Omega_{(\alpha)AB}^* \dot{u}_\alpha^A \dot{u}_\beta^B \cdot n_{(\mu)}^i.$$

We now define secondary normal curvature as follows:

$$(5.12) \quad \frac{1}{R_{(\alpha)}^* R_{(\beta)}^*} \frac{d\Omega}{ds} g_{\alpha\beta}^*(x^h, x_\alpha^h) \frac{D x_\gamma^i}{D s^2} \frac{D x_\delta^j}{D s^2},$$

which on simplification yields:

$$(5.13) \quad \frac{1}{R_{(\alpha)}^* R_{(\beta)}^*} \frac{d\Omega}{ds} = \sum_{\nu} \Omega_{(\alpha)AC}^* \Omega_{(\beta)BD}^* \dot{u}_\alpha^A \dot{u}_\gamma^C \dot{u}_\beta^B \dot{u}_\delta^D,$$

from which we can say that the secondary normal curvature is independent of the choice of secondary normals.

6. COVARIANT DERIVATIVE OF NORMALS. — We define the tensor $n_{(a)}^i$ as the covariant derivative of the vector $n_{(a)}^i$ and it is the projection of $n_{(a)}^i$ onto X_m . Thus

$$(5.1) \quad n_{(a)}^i{}_{;A} = n_{(a)}^i{}_{;k} B_A^k.$$

Since we know that

$$(6.2) \quad n_{(a)}^i{}_{;A} = \frac{\partial n_{(a)}^i}{\partial u^A} - \frac{\partial n_{(a)}^i}{\partial \hat{u}_7^c} \Gamma_{DA}^c \hat{u}_7^D + \Gamma_{ik}^i n_{(a)}^k B_A^k,$$

therefore differentiating relation (3.18) with respect to u^c we get on simplification and rearrangement of terms the following relation

$$(6.3) \quad g_0^{\alpha\beta} B_A^\alpha n_{(a)}^{\beta}{}_{;c} + n_{(a)}^{\beta}{}_{;c} B_A^\alpha \left(\frac{\partial g_0^{\alpha\beta}}{\partial u^c} - g_0^{\alpha\delta} \Gamma_{ik}^{\delta\beta} B_C^k \right) + \frac{\partial n_{(a)}^{\beta}{}_{;c}}{\partial \hat{u}_7^c} \Gamma_{DC}^{\beta} \hat{u}_7^D g_0^{\alpha\beta} B_A^\alpha + g_0^{\alpha\beta} \frac{\partial^2 x^i}{\partial u^A \partial u^c} n_{(a)}^{\beta}{}_{;i} = 0.$$

In equation (6.3) if we add and subtract

$$g_{\beta\alpha}^{\alpha\beta} \Gamma_{ik}^{\beta} B_C^k n_{(a)}^{\alpha}{}_{;c} B_A^\alpha, \quad \frac{\partial g_0^{\alpha\beta}}{\partial \hat{u}_7^c} \Gamma_{DC}^{\beta} \hat{u}_7^D n_{(a)}^{\alpha}{}_{;c},$$

on rearrangement of terms we get

$$(6.4) \quad g_0^{\alpha\beta} n_{(a)}^{\beta}{}_{;c} B_A^\alpha + g_0^{\alpha\beta} \left(\frac{\partial^2 x^i}{\partial u^A \partial u^c} + \Gamma_{ik}^i B_A^k B_C^k \right) n_{(a)}^{\beta}{}_{;i} + \frac{\partial g_0^{\alpha\beta}}{\partial \hat{u}_7^c} \Gamma_{DC}^{\beta} \hat{u}_7^D n_{(a)}^{\alpha}{}_{;c} B_A^\alpha + \frac{\partial n_{(a)}^{\beta}{}_{;c}}{\partial \hat{u}_7^c} \Gamma_{DC}^{\beta} g_0^{\alpha\beta} B_A^\alpha \hat{u}_7^D + \left(\frac{\partial g_0^{\alpha\beta}}{\partial u^c} - g_0^{\alpha\delta} \Gamma_{ik}^{\delta\beta} B_C^k - g_0^{\alpha\delta} B_C^k \Gamma_{ik}^{\delta\beta} - \frac{\partial g_0^{\alpha\beta}}{\partial \hat{u}_7^c} \Gamma_{DC}^{\beta} \hat{u}_7^D \right) \cdot n_{(a)}^{\alpha}{}_{;c} B_A^\alpha = 0.$$

If we put

$$g_{0j|k}^{\alpha\beta} (x^i, x_c^j) \stackrel{\text{def}}{=} B_{ik}^{\alpha\beta} (x^i, x_c^j),$$

in (6.4) we get on simplification

$$(6.5) \quad g_0^{\alpha\beta} B_A^\alpha n_{(a)}^{\beta}{}_{;c} + \frac{\partial g_0^{\alpha\beta}}{\partial \hat{u}_7^c} \Gamma_{DC}^{\beta} \hat{u}_7^D B_A^\alpha n_{(a)}^{\beta}{}_{;c} + \psi_{(a)} \Omega_{0\alpha c} + B_{0ik}^{\alpha\beta} B_A^k n_{(a)}^{\beta}{}_{;i} = 0.$$

Since $n_{(a)}^j c$ is not tangential to X_m we can decompose it as follows :

$$(6.6) \quad n_{(a)}^j c = D_{bc}^E B_K^j + \sum_{\nu} N_{(a)C}^{(\nu)} n_{\nu}^j,$$

where D_{bc}^E and $N_{(a)C}^{(\nu)}$ are to be determined.

Multiplying equation (6.6) by $g_{ij}^{\nu} B_A^i$ and using (6.5) we obtain

$$(6.7) \quad D_{bc}^E = -g_{ab}^{ik} \left[E_{ijk}^{\nu} (x^i, x_{\nu}^j) B_C^k B_A^i n_{\nu}^{*k} + \frac{\partial g_{ij}^{\nu}}{\partial u_{\nu}^l} \Gamma_{bc}^E \hat{u}_{\nu}^D B_A^i n_{\nu}^{*j} + \psi_{(a)} \Omega_{(a)AC}^{*k} \right].$$

If we multiply (6.6) by $n_{(a)}^{*j}$ we obtain

$$(6.8) \quad N_{(a)C}^{(\nu)} \psi_{(\nu)} = n_{(a)}^{*j} c n_{\nu}^{*j}.$$

Now we shall consider the relation

$$(6.9) \quad g_{ij}^{\nu} (x^k, x_{\nu}^l) n_{\nu}^{*i} n_{\nu}^{*j} = \psi_{(\nu)} \delta_{ik}^{\nu},$$

(no summation on μ), if we differentiate (6.9) with respect to u^c and eliminate $\partial n_{\nu}^{*i} / \partial u^c$, from (6.2) and (6.9), we get by virtue of (6.8) on simplification

$$(6.10) \quad \begin{aligned} & \psi_{(\nu)} N_{(a)C}^{(\nu)} + \psi_{(a)} N_{(a)C}^{(\nu)} \\ &= \frac{\partial \psi_{(a)}}{\partial u^c} \delta_{ik}^{\nu} - \frac{\partial \psi_{(a)}}{\partial u_{\nu}^l} \delta_{ik}^{\nu} \Gamma_{bc}^E \hat{u}_{\nu}^D - E_{ijk}^{\nu} n_{\nu}^{*i} n_{\nu}^{*j} B_C^k, \end{aligned}$$

which implies the properties of $N_{(a)C}^{(\nu)}$.

As we have defined $n_{(a)}^j c$, we can also define $n_{(a)}^j c$ by projecting $n_{(a)}^j c$ onto X_m , thus

$$(6.11) \quad n_{(a)}^j c = \frac{\partial n_{(a)}^j}{\partial u^c} - \frac{\partial n_{(a)}^j}{\partial u_{\nu}^l} \Gamma_{bc}^E \hat{u}_{\nu}^D + \Gamma_{bc}^E n_{(a)}^j B_C^k.$$

If we put

$$g_{(a)k}^{\nu} (x^k, n_{\nu}^k) = E_{(a)ijk}^{\nu},$$

we can find

$$(6.12) \quad \Omega_{(a)AC}^{\nu} = -g_{(a)k}^{\nu} B_A^i n_{(a)}^j c - E_{(a)ijk}^{\nu} B_C^k B_A^i n_{\nu}^{*j} - \frac{\partial g_{ij}^{\nu}}{\partial u_{\nu}^l} \Gamma_{bc}^E \hat{u}_{\nu}^D B_A^i n_{\nu}^{*j}.$$

Now we decompose the tensor $n_{(a)}^j c$ as follows :

$$(6.13) \quad n_{(a)}^j c = \Lambda_{bc}^D B_b^j + \sum_{\nu} V_{(a)C}^{(\nu)} n_{\nu}^j,$$

where Λ_{bc}^D and $V_{(a)C}^{(\nu)}$ are to be determined.

Multiplying equation (6.13) by $g_{ij}^{ab}(x^k, n_{\alpha}^b)$ and n_{α}^b respectively we get on simplification

$$(6.14) \quad \Lambda_{\alpha i}^D = -\gamma_{(\alpha)\beta}^{AD} \left\{ \Omega_{\beta i)AC}^A + E_{\beta\alpha\beta i)k}^{A\alpha\beta} B_C^k B_A^i n_{\beta}^i + \frac{\partial g_{ij}^{ab}}{\partial u_{\beta}^b} \Gamma_{BC}^{\alpha\beta} \dot{u}_{\beta}^B B_A^i n_{\alpha}^i \right\},$$

and

$$(6.15) \quad n_{\beta}^i |_{(\alpha)} n_{\alpha}^j = \sum_{\gamma} V_{(\alpha)\gamma}^{(\beta)} \cos(n, n),$$

where

$$\cos(n, n) = n_{(\alpha)}^i n_{(\alpha)}^j.$$

By virtue of equation (6.13) we can write

$$(6.16) \quad \frac{D n_{\alpha}^i}{D s^{\gamma}} = -\gamma_{(\alpha)\beta}^{AD} \left\{ \frac{\partial g_{ij}^{ab}}{\partial u_{\beta}^b} \Gamma_{BC}^{\alpha\beta} B_A^k B_D^i n_{\alpha}^i \dot{u}_{\beta}^B + \Omega_{\beta i)AC}^A B_D^i \dot{u}_{\beta}^C + E_{\beta\alpha\beta i)k}^{A\alpha\beta} B_A^i B_D^k n_{\alpha}^i \dot{x}_{\beta}^k \right\} + \sum_{\gamma} V_{(\alpha)\gamma}^{(\beta)} n_{\beta}^i \dot{u}_{\gamma}^C.$$

Since we know that

$$n_{\alpha}^i = g_{ij}^{ab}(x, n) n_{\alpha}^j,$$

it gives on differentiation

$$(6.17) \quad \frac{D n_{\alpha}^i}{D s^{\gamma}} = g_{ij}^{ab}(x, n) \frac{\partial n_{\alpha}^j}{\partial s^{\gamma}} + E_{\beta\alpha\beta i)k}^{A\alpha\beta} n_{\alpha}^i \dot{x}_{\beta}^k,$$

which by virtue of (6.16) implies:

$$(6.18) \quad \frac{D n_{\alpha}^i}{D s^{\gamma}} = -\gamma_{(\alpha)\beta}^{AD} g_{ij}^{ab}(x, n) \left\{ \Omega_{\beta i)AC}^A \dot{u}_{\beta}^C B_D^i + \frac{\partial g_{ij}^{ab}}{\partial u_{\beta}^b} \Gamma_{BC}^{\alpha\beta} \dot{u}_{\beta}^C \dot{u}_{\beta}^E g_{ij}^{ab} n_{\alpha}^i B_D^k B_A^j \right\} + \sum_{\gamma} V_{(\alpha)\gamma}^{(\beta)} g_{ij}^{ab}(x, n) n_{\alpha}^i \dot{u}_{\beta}^C.$$

Multiplying equation (6.18) by B_{γ}^i we obtain

$$(6.19) \quad \frac{D n_{\alpha}^i}{D s^{\gamma}} B_{\gamma}^i = -\Omega_{\beta i)AC}^A \dot{u}_{\beta}^C - \frac{\partial g_{ij}^{ab}}{\partial u_{\beta}^b} \Gamma_{BC}^{\alpha\beta} \dot{u}_{\beta}^C \dot{u}_{\beta}^E n_{\alpha}^i B_{\gamma}^k + \sum_{\gamma} V_{(\alpha)\gamma}^{(\beta)} g_{ij}^{ab}(x, n) \dot{u}_{\beta}^C n_{\alpha}^i B_{\gamma}^k,$$

which when combined with

$$g_{AB}^{\alpha\beta}(u^C, \dot{u}_{\beta}^E) \dot{u}_{\beta}^E = \Omega_{\beta\alpha AB}^{\alpha\beta} \dot{u}_{\beta}^E R_{(\alpha)}^{\beta}(x^k, \dot{x}_{\beta}^k),$$

yields

$$(6.20) \quad \frac{D n_{\alpha}^i}{D s^{\gamma}} B_{\gamma}^i = -g_{\alpha\beta}^{\alpha\beta}(u^{\beta}, \dot{u}_{\beta}^E) (R_{(\alpha)}^{\beta})^{-1} \dot{u}_{\beta}^C - \frac{\partial g_{ij}^{ab}}{\partial u_{\beta}^b} \Gamma_{BC}^{\alpha\beta} \dot{u}_{\beta}^C \dot{u}_{\beta}^E n_{\alpha}^i B_{\gamma}^k + \sum_{\gamma} V_{(\alpha)\gamma}^{(\beta)} g_{ij}^{ab}(x, n) n_{\alpha}^i \dot{u}_{\beta}^C B_{\gamma}^k.$$

If we choose a particular set of normals n in such a way that the last term of (6.20) vanishes, then (6.20) reduces to ⁽⁶⁾

$$(6.21) \quad \frac{D n_A^i}{D s^r} B_A^i = - g_{AC}^{(6)} (u^D, \hat{u}_E^D) (R_{(6)D}^r)^{-1} \hat{u}_E^C - \frac{\partial g_{mi}^{(6)}}{\partial \hat{u}_E^D} \Gamma_{BC}^E \hat{u}_T^C \hat{u}_T^E n_A^i B_A^m,$$

which is a generalisation of the Rodrigue's formula of subspace of a Finsler space Eliopoulos [5].

7. GAUSS-CODAZZI EQUATIONS. — To obtain the Gauss-Codazzi equations we consider the covariant derivative of B_A^i with respect to u^k , provided we consider the metric of the submanifold X_m , thus

$$(7.1) \quad B_{A|E}^i = \frac{\partial B_A^i}{\partial u^E} - \frac{\partial B_A^i}{\partial \hat{u}_E^D} \Gamma_{DE}^C \hat{u}_A^D - \Gamma_{AE}^D B_D^i,$$

which by virtue of equation (4.5) leads to

$$(7.2) \quad B_{A|C}^i = B_{AC}^i - \Gamma_{ik}^j B_A^k B_C^j.$$

Using equation (5.1) we can write (7.2) as

$$(7.3) \quad B_{A|C}^i = \sum_{\mu} \Omega_{(6)AC}^{*i} (u^D, \hat{u}_E^D) n_A^{\mu} - \Gamma_{ik}^j B_A^k B_C^j,$$

which on further differentiation with respect to u^D and the application of

$$(7.4) \quad B_{A|C|D}^i - B_{A|D|C}^i = K_{ACD}^{*i} B_A^i + B_{A|E}^i T_{CD}^E - \frac{\partial B_A^i}{\partial \hat{u}_E^D} K_{ECD}^E \hat{u}_A^D,$$

implies by virtue of (6.6), (6.7) and (7.2), after a long calculation the following equation

$$(7.5) \quad K_{ACD}^{*i} B_A^i - \frac{\partial B_A^i}{\partial \hat{u}_E^D} K_{ECD}^E \hat{u}_A^D + T_{CD}^E (B_{AE}^i - \Gamma_{ik}^j B_A^k B_E^j) \\ = \left\{ K_{ik}^j B_A^k - \frac{\partial B_A^i}{\partial \hat{u}_E^D} K_{ik}^j \hat{u}_E^D + T_{ik}^m B_m^k (B_{AE}^i - \Gamma_{lm}^j B_A^k B_E^j) \right\} B_C^i B_D^j \\ + \sum_{\mu} n_A^{\mu} (\Omega_{(6)AC|D}^{*\mu} - \Omega_{(6)AD|C}^{*\mu}) + \sum_{\mu} (\Omega_{(6)AC}^{*\mu} \Omega_{(6)ED}^{*\mu} - \Omega_{(6)AD}^{*\mu} \Omega_{(6)EC}^{*\mu}) g_{(6)}^{(6)} B_C^{\mu} B_D^{\mu} \\ - g_{(6)}^{(6)} \frac{\partial g_{im}^{(6)}}{\partial \hat{u}_E^D} (\Omega_{(6)AC}^{*i} \Gamma_{LD}^E - \Omega_{(6)AD}^{*i} \Gamma_{LC}^E) \hat{u}_L^L B_E^i B_C^m n_A^{*m} \\ - E_{ik}^{*j} g_{(6)}^{(6)} \sum_{\mu} (\Omega_{(6)AC}^{*\mu} B_D^k - \Omega_{(6)AD}^{*\mu} B_C^k) n_A^{*i} + \sum_{\mu} \sum_{\nu} (N_{(6)C}^{(6)} \Omega_{(6)AD}^{*\mu} - N_{(6)D}^{(6)} \Omega_{(6)AC}^{*\mu}) n_A^{*i},$$

where K_{ACD}^{*i} , K_{ECD}^E and T_{CD}^E are corresponding terms for the submanifold X_m as already defined by Rund [10] for the manifold X_n .

Multiplying equation (7.5) by $g_0^{\epsilon\theta} B^l$ and $g_0^{\epsilon\theta} n_a^j$ respectively we get on simplification

$$(7.6) \quad g_{EM}^{\epsilon\theta} K_{ACD}^{*E} - g_0^{\epsilon\theta} B_H^l K_{CD}^E \frac{\partial B_A^l}{\partial \dot{x}_A^E} \dot{u}_A^G - T_{CD}^E \Gamma_{jka}^{\epsilon\theta} B_H^l B_A^k B_B^a \\ + E_{jpk}^{\epsilon\theta} \sum_{\mu} (\Omega_{\mu\theta}^{*j} \Gamma_{AC} B_D^k - \Omega_{\mu\theta}^{*j} B_C^k) n_{(A}^{\mu j)} B_H^k \\ + \frac{\partial g_{EM}^{\epsilon\theta}}{\partial \dot{u}_A^E} (\Omega_{\mu\theta}^{*E} \Gamma_{LD}^E - \Omega_{\mu\theta}^{*E} \Gamma_{LD}^E) n_{(A}^{\mu E} B_H^L \dot{u}_L^E \\ = K_{jka}^{\epsilon\theta} B_A^j K_{bc}^k B_D^b B_H^l - g_0^{\epsilon\theta} \frac{\partial B_A^l}{\partial \dot{x}_A^E} K_{jka}^E B_H^l B_C^k B_D^a \dot{x}_A^E + T_{jka}^E B_C^k B_D^b \Gamma_{jka}^{\epsilon\theta} B_H^l B_A^k$$

and

$$(7.7) \quad K_{jka}^{\epsilon\theta} (x^h, \dot{x}_E^h) B_A^h B_C^k B_D^a n_{(A}^{\mu j)} + \psi_{(A} (\Omega_{\mu\theta}^{*a} B - \Omega_{\mu\theta}^{*a} c) \\ - \frac{\partial B_A^l}{\partial \dot{x}_A^E} K_{jka}^E B_C^k B_D^b \dot{x}_A^E n_{(A}^{\mu a} + T_{jka}^E B_C^k B_D^b (\Omega_{\mu\theta}^{*a} \psi_{(A} B_n^k - \Gamma_{jka}^E B_A^k n_{(A}^{\mu a)}) \\ - E_{jpk}^{\epsilon\theta} \sum_{\mu} (\Omega_{\mu\theta}^{*a} B_B^k - \Omega_{\mu\theta}^{*a} B_C^k) n_{(A}^{\mu a} n_{(A}^{\mu j)} \\ + \sum_{\mu} \psi_{(A} (N_{\mu\theta}^{\epsilon\theta} \Omega_{\mu\theta}^{*a} - N_{\mu\theta}^{\epsilon\theta} \Omega_{\mu\theta}^{*a}) - T_{CD}^E (\psi_{(A} \Omega_{\mu\theta}^{*a} - \Gamma_{jka}^E B_A^k B_B^a n_{(A}^{\mu j)}) \\ + \frac{\partial B_A^l}{\partial \dot{u}_A^E} K_{GCD}^E n_{(A}^{\mu a} \dot{u}_A^G = 0,$$

where

$$\Gamma_{jka}^{\epsilon\theta} = g_0^{\epsilon\theta} \Gamma_{jk}^a \text{ and } K_{jka}^{\epsilon\theta} = g_0^{\epsilon\theta} K_{jka}^a.$$

Equations (7.6) and (7.7) are the Gauss-Codazzi equations for the normals n_a^j in a submanifold X_m of a manifold X_n .

To obtain the Gauss-Codazzi equations for the normals n_a^i , we decompose B_A^k (considered as a vector with respect to the index i) into components along the tangent plane at the point considered and normals n_a^i . Thus we have

$$(7.8) \quad B_{CD}^i = \sum_{\mu} A_{\mu\theta}^i B_{CD}^{\mu} + W_{CD}^i,$$

where $A_{\mu\theta}^i$ and W_{CD}^i are to be determined and W_{CD}^i satisfies

$$W_{CD}^i n_{(A}^{\mu a} = 0.$$

Multiplying equation (7.8) by $n_{(A}^{\mu a}$ we obtain

$$(7.9) \quad \Omega_{\mu\theta}^a = n_{(A}^{\mu a} B_{CD}^i = \sum_{\mu} A_{\mu\theta}^a \cos \left(\frac{n \cdot n}{\theta} \right),$$

which implies

$$(7.10) \quad W_{CD}^i = \sum_{\mu} \Omega_{\mu C D}^{\alpha} n_{\alpha}^i - \sum_{\mu} A_{\mu C D}^{\alpha} n_{\alpha}^i.$$

Using equation (7.2) in (7.10) we get

$$(7.11) \quad B_{CD}^i = \sum_{\mu} A_{\mu C D}^{\alpha} n_{\alpha}^i + W_{CD}^i - \Gamma_{hk}^i B_{\mu}^h B_{\nu}^k.$$

In analogy with the Finsler space Eliopoulos [5] we can also express W_{CD}^i as follows:

$$(7.12) \quad W_{CD}^i = B_{\mu}^i \sum_{\nu} A_{\nu C D}^{\alpha} M_{\alpha}^{\nu},$$

where

$$M_{\alpha}^{\nu} = n_{\alpha}^i B_{\nu}^i.$$

From (7.12) we can easily obtain

$$(7.13) \quad W_{CD}^i = B_{\mu}^i \left(\sum_{\nu} A_{\nu C D}^{\alpha} M_{\alpha}^{\nu} + \sum_{\nu} A_{\nu C D}^{\alpha} M_{\alpha}^{\nu} + \sum_{\nu} (A_{\nu C D}^{\alpha} M_{\alpha}^{\nu}) B_{\mu}^i \right).$$

Now differentiating (7.11) with respect to the metric of the submanifold X_m , interchanging the last two indices, subtracting the resulting equation from first and using equations (6.11), (7.4) and (7.13) we get on simplification the following relation:

$$(7.14) \quad K_{CD}^E B_{\mu}^i - \frac{\partial B_{\mu}^i}{\partial \dot{x}_{\alpha}^E} K_{CD}^{\alpha} \dot{x}_{\alpha}^E + T_{DF}^E (B_{CE}^i - \Gamma_{hk}^i B_{\mu}^h B_{\nu}^k) \\ = \left\{ K_{hk}^i B_{\mu}^h - \frac{\partial B_{\mu}^i}{\partial \dot{x}_{\alpha}^E} K_{\mu\alpha}^E \dot{x}_{\alpha}^E + T_{DF}^E B_{\mu}^E (B_{CE}^i - \Gamma_{hk}^i B_{\mu}^h B_{\nu}^k) \right\} B_{\nu}^j B_{\rho}^k \\ + B_{\mu}^i \left[\sum_{\nu} \gamma_{\nu\alpha\beta}^{jk} (\Omega_{\nu\alpha\beta}^k A_{\mu C D}^{\alpha} - \Omega_{\nu\alpha\beta}^k A_{\mu C D}^{\alpha}) \right] \\ + \sum_{\nu} M_{\alpha}^{\nu} (A_{\nu C D}^{\alpha} B_{\mu}^i - A_{\mu C D}^{\alpha} B_{\nu}^i) + \sum_{\nu} (M_{\alpha}^{\nu} A_{\nu C D}^{\alpha} - M_{\alpha}^{\nu} A_{\mu C D}^{\alpha}) \\ - \sum_{\nu} g_{\alpha\beta}^{\gamma\delta} E_{\nu\beta\alpha k}^{\gamma} n_{\alpha}^k (A_{\nu C D}^{\alpha} B_{\mu}^i - A_{\mu C D}^{\alpha} B_{\nu}^i) + \sum_{\nu} n_{\alpha}^i (A_{\nu C D}^{\alpha} B_{\mu}^i - A_{\mu C D}^{\alpha} B_{\nu}^i) \\ + \sum_{\nu} \sum_{\rho} (A_{\nu C D}^{\alpha} \Omega_{\rho C D}^{\alpha} - A_{\mu C D}^{\alpha} \Omega_{\rho C D}^{\alpha}) n_{\alpha}^i M_{\alpha}^{\nu} + \sum_{\nu} \sum_{\rho} n_{\alpha}^i (A_{\nu C D}^{\alpha} V_{\rho}^{\beta} - A_{\mu C D}^{\alpha} V_{\rho}^{\beta}).$$

Multiplying equation (7.14) by $\delta_{\mu}^{e0} B_{\mu}^i$ and solving we get

$$(7.15) \quad \delta_{\mu}^{e0} \gamma_{\nu\alpha\beta}^{jk} [K_{CD}^E - \sum_{\rho} \gamma_{\rho\alpha\beta}^{jk} (\Omega_{\nu\alpha\beta}^k A_{\mu C D}^{\alpha} - \Omega_{\nu\alpha\beta}^k A_{\mu C D}^{\alpha}) \\ - \sum_{\rho} M_{\alpha}^{\nu} (A_{\nu C D}^{\alpha} B_{\mu}^i - A_{\mu C D}^{\alpha} B_{\nu}^i) - \sum_{\rho} (M_{\alpha}^{\nu} A_{\nu C D}^{\alpha} - M_{\alpha}^{\nu} A_{\mu C D}^{\alpha})]$$

$$\begin{aligned}
 & - \frac{\Delta B_i^c}{\Delta \hat{u}_a^A} K_{0DF}^A \hat{u}_a^c \hat{g}_0^{c0} B_H^i + T_{DF}^{c0} \left[\sum_{\mu} n_a^i (A_{DQCE}^{\mu} \hat{g}_0^{c0} B_H^i + g_{LH}^{c0} B_L^i A_{DQCE}^{\mu}) \right. \\
 & \left. - \Gamma_{\mu a}^{c0} B_C^{\mu} B_E^a B_H^i \right] + \sum_{\mu} g_{00}^{\mu} E_{(a)ak}^{(0)\mu} \hat{g}_0^{c0} (x, n) n_a^i B_H^i (A_{DQCD}^{\mu} B_F^k - A_{DQCF}^{\mu} B_D^k) \\
 & \quad - \sum_{\mu} \hat{g}_0^{c0} (x, n) n_a^i B_H^i (A_{DQCD}^{\mu} \Gamma - A_{DQCF}^{\mu} D) \\
 & \quad - \sum_{\mu} \sum_{\nu} (A_{DQCD}^{\mu} V_{(a)F}^{(\nu)\mu} - A_{DQCF}^{\mu} V_{(a)D}^{(\nu)\mu}) \hat{g}_0^{c0} n_a^i B_H^i \\
 = & \left[K_{\mu a}^{c0} B_H^i B_C^{\mu} - \frac{\Delta B_i^c}{\Delta \hat{x}_a^m} K_{\mu a}^m \hat{x}_a^c \hat{g}_0^{c0} B_H^i + T_{\mu a}^{c0} B_E^a \left\{ \sum_{\mu} A_{DQCE}^{\mu} n_a^i (\hat{g}_0^{c0} B_H^i + g_{LH}^{c0} B_L^i) \right. \right. \\
 & \quad \left. \left. - \Gamma_{\mu a}^{c0} B_C^{\mu} B_E^a B_H^i \right\} \right] B_D^k B_F^k.
 \end{aligned}$$

Similarly multiplying equation (7.14) by $\hat{g}_0^{c0} (x, n) n_a^i$ we get

$$\begin{aligned}
 (7.16) \quad & \left\{ K_{\mu a}^{c0} B_C^{\mu} - \frac{\Delta B_i^c}{\Delta \hat{x}_a^m} K_{\mu a}^m \hat{g}_0^{c0} \hat{x}_a^c \right\} n_a^i B_D^k B_F^k + \frac{\Delta B_i^c}{\Delta \hat{u}_a^A} K_{0DF}^A \hat{u}_a^c \hat{g}_0^{c0} (x, n) n_a^i \\
 & - T_{DF}^{c0} \left(\sum_{\mu} \Omega_{DQCE}^{\mu} \hat{g}_0^{c0} n_a^i n_b^j - \Gamma_{\mu a}^{c0} B_C^{\mu} B_E^a n_b^j \right) \\
 & + T_{\mu a}^{c0} \hat{g}_0^{c0} (x, n) B_E^a n_b^j B_D^k B_F^k \left(\sum_{\mu} \Omega_{DQCE}^{\mu} n_a^i - \Gamma_{\mu a}^{c0} B_C^{\mu} B_E^a \right) \\
 & - \sum_{\mu} g_{00}^{\mu} E_{DQak}^{(0)\mu} \hat{g}_0^{c0} (x, n) (A_{DQCD}^{\mu} B_F^k - A_{DQCF}^{\mu} B_D^k) n_a^i n_b^j \\
 & \quad + \sum_{\mu} \hat{g}_0^{c0} (x, n) n_a^i n_b^j (A_{DQCD}^{\mu} \Gamma - A_{DQCF}^{\mu} D) \\
 & + \sum_{\mu} \sum_{\nu} M_{\mu\nu}^k \hat{g}_0^{c0} (x, n) n_a^i n_b^j (A_{DQCD}^{\mu} \Omega_{(a)EF}^{\nu} - A_{DQCF}^{\mu} \Omega_{(a)ED}^{\nu}) + \\
 & + \sum_{\mu} \sum_{\nu} \hat{g}_0^{c0} (x, n) n_a^i n_b^j (A_{DQCD}^{\mu} V_{(a)F}^{(\nu)\mu} - A_{DQCF}^{\mu} V_{(a)D}^{(\nu)\mu}) = 0.
 \end{aligned}$$

Equations (7.15) and (7.16) are Gauss-Codazzi equations for the normals n_a^i in the submanifold X_m of the manifold X_n .

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