

On areal metrics on complex manifolds (**)

Abstract: The 4-index areal metric tensor has been defined and its properties have been studied by Rund [3], [4]. The geometry of areal spaces in these works are based on real coordinate system. The purpose of this paper is to investigate a certain class of 4-index metric tensor on complex manifolds. After giving the outline of complex manifolds in §1, we introduce this 4-index metric tensor in §2. The section 3 is devoted to a fairly detailed description of the many identities which result from the homogeneity condition of the metric function. In §4 these preliminary results, all of which are basically concerned with a 4-index metric tensor, are used in the construction of suitable connection coefficients, which are explicitly derivable from the metric tensor and its various derivatives. The last section is devoted in construction of the covariant partial derivatives of a given vector field with the help of these connection coefficients. Our considerations are purely local in character. In some places the detailed calculations have been suppressed for the sake of brevity. With regard to such instances reference is made to the report [3]. Throughout this paper the Latin indices i, j, h, k, l, m, p, q run over 1 to a while Greek indices $\alpha, \beta, \gamma, \lambda$ run over 1 to m .

1. - INTRODUCTION

We consider a $2n$ -dimensional real manifold X_{2n} (of class C^∞) referred to local coordinates (x^i, y^j) . Corresponding to each point P of X_{2n} we introduce complex numbers z^i ,

$$(1.1) \quad z^i = x^i + i y^i \quad (i^2 = -1)$$

which may be regarded as the complex coordinate of P (with respect to the given coordinate system). If there exist complex coordinate neighbourhoods $U(z^i)$, $U(\bar{z}^i)$ (where \bar{z}^i refer to another local coordinate system) such that in the intersection of these neighbourhoods, we have

$$(1.2) \quad \bar{z}^i = \bar{z}^i(z^h) \quad \det \begin{bmatrix} \frac{\partial \bar{z}^i}{\partial z^h} \end{bmatrix} \neq 0$$

where $\bar{z}^i(z^h)$ are holomorphic functions of z^h , then space X_{2n} is said to admit a complex structure. Under these circumstances X_{2n} is called a complex space of complex dimension n and is denoted by C_n .

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With (1.1) we may associate the conjugate complex

$$(1.3) \quad \bar{z}^j = x^j - i y^j$$

so that (1.2) carries with it the corresponding conjugate complex transformation,

$$(1.4) \quad \bar{z}^{j*} = \bar{z}^j (z^{j*}).$$

An analytic m -dimensional subspace C_m of C_n ($m < n$) is represented parametrically by equations ([1] page 104)

$$(1.5) \quad z^j = z^j(u^a), \quad z^{j*} = z^{j*}(u^{a*})$$

in which the z^j, z^{j*} are holomorphic functions of the complex variables u^a, u^{a*} respectively. Thus the derivatives $z_a^j = \frac{\partial z^j}{\partial u^a}$ and their complex conjugate $z_a^{j*} = \overline{\frac{\partial z^j}{\partial u^a}}$ are defined, each of which are the elements of an $n \times m$ matrix which is always supposed to be of rank m .

2. - FUNDAMENTAL FORMULAE

We consider real Lagrange function L of the form,

$$(2.1) \quad L = L(z^j, z^{j*}, \dot{z}_a^j, \dot{z}_a^{j*})$$

satisfying the conditions

(A) The Lagrangian L is of class C^4 in all its arguments and it is scalar with respect to transformations (1.2) and (1.4).

(B) The Lagrangian L is positive for all independent sets of arguments $\dot{z}_a^j, \dot{z}_a^{j*}$.

(C) The integral

$$(2.2) \quad I = \int_G L du^1 \wedge \dots \wedge du^m \wedge du^{1*} \wedge \dots \wedge du^{m*}$$

over a fixed region G of C_m is invariant under the holomorphic transformations of the complex parameters

$$(2.3) \quad \bar{u}^a = \bar{u}^a(u^b), \quad \bar{u}^{a*} = \bar{u}^{a*}(u^{b*})$$

(D) The $n \times m$ determinant

$$D = \det \begin{bmatrix} \frac{2m}{2} & \frac{\partial^2 L}{\partial \dot{z}_a^j \partial \dot{z}_a^{j*}} \end{bmatrix}$$

is non vanishing for linearly independent $\dot{z}_a^j, \dot{z}_a^{j*}$.

The condition C is equivalent to the relations [2]

$$(2.4) \quad (a) \frac{\partial L}{\partial z^i_\alpha} \dot{z}^i_\alpha = \delta^i_\alpha L \quad (b) \frac{\partial L}{\partial z^{i*}_\alpha} \dot{z}^{i*}_\alpha = \delta^{i*}_\alpha L.$$

In particular if we put $\alpha = \beta$ in (2.4) we get

$$(2.5) \quad (a) \frac{\partial L}{\partial z^i_\alpha} \dot{z}^i_\alpha = m L \quad ; \quad (b) \frac{\partial L}{\partial z^{i*}_\alpha} \dot{z}^{i*}_\alpha = m L.$$

Differentiating both the equations of (2.5) with respect to \dot{z}^i_α and adding we get

$$(2.6) \quad \frac{\partial^2 L}{\partial \dot{z}^i_\alpha \partial z^i_\alpha} \dot{z}^i_\alpha + \frac{\partial^2 L}{\partial \dot{z}^i_\alpha \partial z^{i*}_\alpha} \dot{z}^{i*}_\alpha = (2m - 1) \frac{\partial L}{\partial \dot{z}^i_\alpha}.$$

From (2.5) and (2.6) one can find

$$(2.7) \quad \frac{2m}{2} \left[\frac{\partial^2 L}{\partial \dot{z}^i_\alpha \partial z^i_\alpha} \dot{z}^i_\alpha + \frac{\partial^2 L}{\partial \dot{z}^i_\alpha \partial z^{i*}_\alpha} \dot{z}^{i*}_\alpha \right] = L^{\frac{2}{2m-1}} \frac{\partial L}{\partial \dot{z}^i_\alpha}$$

alongwith

$$(2.8) \quad \frac{2m}{2} \left[\frac{\partial^2 L}{\partial \dot{z}^{i*}_\alpha \partial z^i_\alpha} \dot{z}^i_\alpha + \frac{\partial^2 L}{\partial \dot{z}^{i*}_\alpha \partial z^{i*}_\alpha} \dot{z}^{i*}_\alpha \right] = L^{\frac{2}{2m-1}} \frac{\partial L}{\partial \dot{z}^{i*}_\alpha}.$$

Multiplying (2.7) by \dot{z}^i_α and (2.8) by \dot{z}^{i*}_α and adding we get (in view of (2.5))

$$(2.9) \quad 2m L^{\frac{2}{2m}} = g^i_\alpha \dot{z}^i_\alpha \dot{z}^i_\alpha + 2g^{i*}_\alpha \dot{z}^i_\alpha \dot{z}^{i*}_\alpha + g^{i*}_\alpha \dot{z}^{i*}_\alpha \dot{z}^{i*}_\alpha$$

where

$$(2.10) \quad g^i_\alpha (z^i, z^{i*}, \dot{z}^i_\alpha, \dot{z}^{i*}_\alpha) = \frac{2m}{2} \frac{\partial^2 L}{\partial \dot{z}^i_\alpha \partial z^i_\alpha},$$

$$(2.11) \quad g^{i*}_\alpha (z^i, z^{i*}, \dot{z}^i_\alpha, \dot{z}^{i*}_\alpha) = \frac{2m}{2} \frac{\partial^2 L}{\partial \dot{z}^i_\alpha \partial z^{i*}_\alpha},$$

$$(2.12) \quad g^{i*}_\alpha (z^i, z^{i*}, \dot{z}^i_\alpha, \dot{z}^{i*}_\alpha) = \frac{2m}{2} \frac{\partial^2 L}{\partial \dot{z}^{i*}_\alpha \partial z^{i*}_\alpha}.$$

From (2.9) it is evident that if L is interpreted as measure of the area dA of an m-dimensional complex subspace (2m-dimensional real subspace) spanned by $\dot{z}^i_\alpha \dot{z}^{i*}_\alpha$ at the points z^i, z^{i*} of C_m in the sense that

$$(2.13) \quad dA = L(z^i, z^{i*}, \dot{z}^i_\alpha, \dot{z}^{i*}_\alpha) du^1 \wedge \dots \wedge du^m \wedge du^{1*} \wedge \dots \wedge du^{m*}$$

then the tensors (2.10), (2.11) and (2.12) can be regarded as a suitable areal metric tensor ([3] page 289).

It is to be noted that g_{ij}^{st} is symmetric in pairs of indices such as (s, h), (r, j). The similar symmetries exist for the tensors g_{ij}^{rst} and g_{ij}^{rstu} . Furthermore $g_{ij}^{rst} \neq g_{ij}^{rst}$.

3. IDENTITIES RESULTING FROM HOMOGENEITY CONDITION (2.4)

Differentiating (2.5) with respect to z_3^h and $z_{3^*}^{h^*}$ respectively we get after some simplifications

$$(3.1) \quad (a) \quad \frac{\partial^3 L}{\partial z_3^h \partial z_4^i} z_4^i = (m-1) \frac{\partial L}{\partial z_3^h} \quad (b) \quad \frac{\partial^3 L}{\partial z_3^h \partial z_{4^*}^{i^*}} z_{4^*}^{i^*} = m \frac{\partial L}{\partial z_3^h}$$

$$(3.2) \quad (a) \quad \frac{\partial^3 L}{\partial z_{3^*}^{h^*} \partial z_4^i} z_4^i = m \frac{\partial L}{\partial z_{3^*}^{h^*}} \quad (b) \quad \frac{\partial^3 L}{\partial z_{3^*}^{h^*} \partial z_{4^*}^{i^*}} z_{4^*}^{i^*} = (m-1) \frac{\partial L}{\partial z_{3^*}^{h^*}}$$

But we have from (2.10) and (2.11)

$$(3.3) \quad g_{ij}^{st} = \left(\frac{2}{2m} - 1 \right) L^{\frac{2}{2m}-2} \frac{\partial L}{\partial z_3^h} \frac{\partial L}{\partial z_4^i} + L^{\frac{2}{2m}-1} \frac{\partial^2 L}{\partial z_3^h \partial z_4^i}$$

$$(3.4) \quad g_{ij}^{s^*t^*} = \left(\frac{2}{2m} - 1 \right) L^{\frac{2}{2m}-2} \frac{\partial L}{\partial z_3^h} \frac{\partial L}{\partial z_{4^*}^{i^*}} + L^{\frac{2}{2m}-1} \frac{\partial^2 L}{\partial z_3^h \partial z_{4^*}^{i^*}}$$

Differentiating (3.4) with respect to z_7^k we get,

$$(3.5) \quad \frac{\partial g_{ij}^{s^*t^*}}{\partial z_7^k} = \left(\frac{2}{2m} - 1 \right) \left\{ \left(\frac{2}{2m} - 2 \right) L^{\frac{2}{2m}-3} \frac{\partial L}{\partial z_3^h} \frac{\partial L}{\partial z_4^i} \frac{\partial L}{\partial z_{4^*}^{i^*}} \right. \\ \left. + L^{\frac{2}{2m}-2} \left(\frac{\partial^3 L}{\partial z_7^k \partial z_3^h \partial z_4^i} \frac{\partial L}{\partial z_{4^*}^{i^*}} + \frac{\partial L}{\partial z_3^h} \frac{\partial^2 L}{\partial z_7^k \partial z_{4^*}^{i^*}} \right) \right. \\ \left. + L^{\frac{2}{2m}-2} \frac{\partial L}{\partial z_7^k} \frac{\partial^2 L}{\partial z_3^h \partial z_{4^*}^{i^*}} \right\} + L^{\frac{2}{2m}-1} \frac{\partial^3 L}{\partial z_7^k \partial z_3^h \partial z_{4^*}^{i^*}}$$

Multiplying (3.5) by z_7^k and using (2.5), (3.1) a and (3.2) a we get after some simplifications,

$$(3.6) \quad \frac{\partial g_{ij}^{s^*t^*}}{\partial z_7^k} z_7^k = L^{\frac{2}{2m}-1} \left\{ (1-m) \frac{\partial^3 L}{\partial z_3^h \partial z_{4^*}^{i^*}} + \frac{\partial^3 L}{\partial z_7^k \partial z_3^h \partial z_{4^*}^{i^*}} z_7^k \right\}.$$

However differentiation of (3.1) a with respect to $z_{7^*}^{k^*}$ gives

$$(3.7) \quad \frac{\partial^3 L}{\partial z_{7^*}^{k^*} \partial z_3^h \partial z_4^i} z_4^i = (m-1) \frac{\partial^2 L}{\partial z_{7^*}^{k^*} \partial z_3^h}$$

and this shows that (3.6) reduces to

$$(3.8) \quad \frac{\partial g_{ij}^{h*}}{\partial z_j^k} z_j^k = 0.$$

Since $\frac{\partial g_{ij}^{h*}}{\partial z_j^k}$ is symmetric in the pairs of indices (j, h); (i, k) it also follows from (3.8) that

$$(3.9) \quad \frac{\partial g_{ij}^{h*}}{\partial z_j^h} z_j^h = 0.$$

Similarly we can show that

$$(3.10) \quad \frac{\partial g_{ij}^{h*}}{\partial z_i^k} z_i^k = 0 = \frac{\partial g_{ij}^{h*}}{\partial z_i^h} z_i^h.$$

On the other hand multiplying (3.5) by z_i^h and using (2.5) and (3.1) b we obtain

$$(3.11) \quad \frac{\partial g_{ij}^{h*}}{\partial z_j^k} z_i^h = \left(\frac{2}{2m} - 1 \right) L_i^{\frac{2}{2m} - 2} \frac{\partial L}{\partial z_j^k} \frac{\partial L}{\partial z_i^h} + \\ + L_i^{\frac{2}{2m} - 1} \left\{ \frac{\partial^2 L}{\partial z_j^k \partial z_i^h} (1-m) + \frac{\partial^2 L}{\partial z_j^k \partial z_i^h \partial z_i^h} z_i^h \right\}.$$

Differentiation of (3.1) b with respect to z_j^k yields

$$(3.12) \quad \frac{\partial^2 L}{\partial z_j^k \partial z_i^h \partial z_i^h} z_i^h = m \frac{\partial^2 L}{\partial z_j^k \partial z_i^h}.$$

Applying (3.12) to (3.11) we find that

$$\frac{\partial g_{ij}^{h*}}{\partial z_j^k} z_i^h = \left(\frac{2}{2m} - 1 \right) L_i^{\frac{2}{2m} - 2} \frac{\partial L}{\partial z_j^k} \frac{\partial L}{\partial z_i^h} + L_i^{\frac{2}{2m} - 1} \frac{\partial^2 L}{\partial z_j^k \partial z_i^h}$$

which in view of (3.3) yields

$$(3.13) \quad \frac{\partial g_{ij}^{h*}}{\partial z_j^k} z_i^h = g_{ik}^{hj}.$$

Similarly we can show that

$$(3.14) \quad \frac{\partial g_{ij}^{h*}}{\partial z_i^k} z_i^h = g_{ik}^{hj*}.$$

Multiplying (3.13) by z_j^h and using (3.9) we find

$$(3.15) \quad g_{ik}^{hj} z_j^h = 0$$

and similarly

$$(3.16) \quad g_{ik}^{hj*} z_j^h = 0.$$

Furthermore from (3.3), (3.1) a and (2.5) it follows that

$$(3.17) \quad \frac{\partial g_{ij}^{\alpha}}{\partial z_{\gamma}^k} z_{\gamma}^k = - \left(\frac{2}{2m} - 1 \right) L^{\frac{2}{2m} - 1} \frac{\partial L}{\partial z_{\beta}^h} \frac{\partial L}{\partial z_{\alpha}^l} + \\ + L^{\frac{2}{2m} - 1} \left\{ (1-m) \frac{\partial^2 L}{\partial z_{\gamma}^k \partial z_{\alpha}^l} + \frac{\partial^2 L}{\partial z_{\gamma}^k \partial z_{\beta}^h \partial z_{\alpha}^l} z_{\gamma}^k \right\}$$

while differentiation of (3.1) a gives

$$(3.18) \quad \frac{\partial^2 L}{\partial z_{\gamma}^k \partial z_{\beta}^h \partial z_{\alpha}^l} z_{\gamma}^k = (m-2) \frac{\partial^2 L}{\partial z_{\beta}^h \partial z_{\alpha}^l}.$$

Substituting (3.18) in (3.17) and using (3.3) once more, we obtain,

$$(3.19) \quad \frac{\partial g_{ij}^{\alpha}}{\partial z_{\gamma}^k} z_{\gamma}^k = -g_{ij}^{\alpha}$$

and by symmetry in the pairs of indices (j, h), (γ, k) we get

$$(3.20) \quad \frac{\partial g_{ij}^{\alpha}}{\partial z_{\beta}^h} z_{\gamma}^k = -g_{ij}^{\alpha}.$$

Similarly it can be shown that

$$(3.21) \quad \frac{\partial g_{i\alpha}^{\beta\gamma}}{\partial z_{\gamma}^k} z_{\gamma}^k = \frac{\partial g_{i\alpha}^{\beta\gamma}}{\partial z_{\beta}^h} z_{\gamma}^k = -g_{i\alpha}^{\beta\gamma}.$$

A direct application of relations (3.15), (3.16), (3.19) and (3.21) will yield,

$$(3.22) \quad \frac{\partial g_{ij}^{\alpha}}{\partial z_{\gamma}^k} z_{\gamma}^k z_{\beta}^h = 0 \quad ; \quad \frac{\partial g_{i\alpha}^{\beta\gamma}}{\partial z_{\gamma}^k} z_{\gamma}^k z_{\beta}^h = 0.$$

Also in consequence of (3.15) and (3.16) the relation (2.9) reduces to

$$(3.23) \quad m L^{\frac{1}{m}} = g_{i\alpha}^{\beta\gamma} z_{\beta}^h z_{\gamma}^k.$$

4. - THE CONNECTION COEFFICIENTS

Let us suppose that we are given an m-dimensional complex subspace C_m . Under the holomorphic transformations of the form (1.2) and (1.4) of the local coordinates of C_m , the quantities z_{α}^i and z_{α}^{i*} transform as follows,

$$(4.1) \quad z_{\alpha}^i = B_{\alpha}^i(\bar{z}^j) z_{\alpha}^h \quad ; \quad z_{\alpha}^{i*} = B_{\alpha}^{i*}(\bar{z}^{j*}) z_{\alpha}^{h*}$$

where we have written

$$(4.2) \quad B_{\alpha}^i(\bar{z}^j) = \frac{\partial z^i}{\partial z^h} \quad ; \quad B_{\alpha}^{i*}(\bar{z}^{j*}) = \frac{\partial z^{i*}}{\partial z^{j*}}.$$

We shall also put

$$(4.3) \quad B_{ia}^j = \frac{\partial^2 z^j}{\partial \bar{z}^i \partial \bar{z}^a} = B_{ih}^j, \quad B_{i\alpha\beta}^{\mu\nu} = \frac{\partial^2 z^{\mu\nu}}{\partial \bar{z}^i \partial \bar{z}^{\alpha\beta}} = B_{i\alpha\beta}^{\mu\nu}.$$

Since the functions (1.2) and (1.4) are assumed to be holomorphic, we have,

$$(4.4) \quad \frac{\partial z^j}{\partial \bar{z}^i \alpha} = 0; \quad \frac{\partial z^{\mu\nu}}{\partial \bar{z}^i} = 0$$

and

$$(4.5) \quad \frac{\partial B_{ia}^j}{\partial \bar{z}^k \alpha} = 0; \quad \frac{\partial B_{i\alpha\beta}^{\mu\nu}}{\partial \bar{z}^k} = 0.$$

The equations (4.2), (4.3), (4.4) and (4.5) have been frequently used in the remaining part of this paper.

For future reference we note that

$$(4.6) \quad \frac{\partial \dot{z}_a^j}{\partial u^i} = \dot{z}_{ia}^j = B_{ih}^j \dot{z}_a^h + B_{i\alpha}^j \dot{z}_a^\alpha,$$

$$(4.7) \quad \frac{\partial \dot{z}_{\alpha\beta}^{\mu\nu}}{\partial u^i} = \dot{z}_{i\alpha\beta}^{\mu\nu} = B_{ih}^{\mu\nu} \dot{z}_{\alpha\beta}^h + B_{i\alpha\beta}^{\mu\nu} \dot{z}_{\alpha\beta}^{\mu\nu}.$$

Clearly, the quantities \dot{z}_a^j and $\dot{z}_{\alpha\beta}^{\mu\nu}$ do not represent components of a tensor field and consequently we must construct the corresponding covariant derivatives.

Since the given Lagrangian L is supposed to be scalar with respect to the transformations (1.2) and (1.4) it follows directly from (4.1) that the quantities (2.10), (2.11) and (2.12) represent components of a covariant tensor of rank two. We, therefore, have,

$$(4.8) \quad \bar{g}_{ij}^{\alpha\beta}(\bar{z}^i, \bar{z}^j, \dot{z}_a^i, \dot{z}_{\alpha\beta}^i) = g_{km}^{\alpha\beta}(z^i, z^j, \dot{z}_a^i, \dot{z}_{\alpha\beta}^i) B_h^k B_p^m.$$

Differentiating (4.8) with respect to \bar{z}^i we get

$$(4.9) \quad \frac{\partial \bar{g}_{ij}^{\alpha\beta}}{\partial \bar{z}^i} = \frac{\partial g_{km}^{\alpha\beta}}{\partial z^i} B_h^k B_p^m + \\ + \frac{\partial g_{km}^{\alpha\beta}}{\partial \dot{z}_a^i} B_{i\alpha}^k \dot{z}_a^m B_h^k B_p^m + g_{km}^{\alpha\beta} B_h^k B_p^m \dot{z}_a^i.$$

Multiplying (4.9) by \dot{z}_a^i and using (4.1), (3.10) we find,

$$(4.10) \quad \frac{\partial \bar{g}_{ij}^{\alpha\beta}}{\partial \bar{z}^i} \dot{z}_a^i = \left\{ \frac{\partial g_{km}^{\alpha\beta}}{\partial z^i} B_h^k \dot{z}_a^m + g_{km}^{\alpha\beta} B_{i\alpha}^k \dot{z}_a^i \right\} B_h^k.$$

If, in addition, this is multiplied by \dot{z}_a^i , we find

$$(4.11) \quad \frac{\partial \bar{g}_{ij}^{\alpha\beta}}{\partial \bar{z}^i} \dot{z}_a^i \dot{z}_a^i = B_h^k \left\{ \frac{\partial g_{km}^{\alpha\beta}}{\partial z^i} \dot{z}_a^i \dot{z}_a^m + g_{km}^{\alpha\beta} B_{i\alpha}^k \dot{z}_a^i \dot{z}_a^i \right\}.$$

Now multiplying (4.9) by \hat{z}_5^k and using (4.1), (3.14) we get

$$(4.12) \quad \frac{\partial \hat{g}_{kl}^{j*}}{\partial \hat{z}^l} \hat{z}_5^k = \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^m} B_l^i B_m^j \hat{z}_5^k + g_{km}^{j*} B_l^i B_m^j \hat{z}_5^k B_l^m \\ + g_{km}^{j*} \hat{z}_5^k B_l^m B_l^j.$$

Again differentiating (4.8) with respect to \hat{z}^l we get

$$(4.13) \quad \frac{\partial \hat{g}_{kl}^{j*}}{\partial \hat{z}^l} = \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^m} B_l^i B_m^j B_l^k + \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^k} B_l^i \hat{z}_5^k B_h^j B_l^m \\ + g_{km}^{j*} B_h^k B_l^j B_l^m.$$

Multiplication of (4.13) by \hat{z}_5^m yields (in view of (4.1) and (3.13)),

$$(4.14) \quad \frac{\partial \hat{g}_{kl}^{j*}}{\partial \hat{z}^l} \hat{z}_5^m = \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^m} B_l^i B_h^k \hat{z}_5^m + g_{kl}^{j*} B_l^i \hat{z}_5^k B_h^j B_l^m \\ + g_{km}^{j*} \hat{z}_5^m B_h^k B_l^j.$$

Further multiplication of (4.14) by \hat{z}_7^l gives

$$(4.15) \quad \frac{\partial \hat{g}_{kl}^{j*}}{\partial \hat{z}^l} \hat{z}_5^m \hat{z}_7^l = \left\{ \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^m} \hat{z}_7^l \hat{z}_5^m + g_{kl}^{j*} B_l^i \hat{z}_5^k \hat{z}_7^l \right\} B_h^j B_l^m \\ + g_{km}^{j*} \hat{z}_5^m B_h^k B_l^j \hat{z}_7^l.$$

Now let us interchange l and h in (4.13) and subtract the result from (4.13). We thus obtain,

$$(4.16) \quad \frac{\partial \hat{g}_{kl}^{j*}}{\partial \hat{z}^l} - \frac{\partial \hat{g}_{lh}^{j*}}{\partial \hat{z}^h} = \left\{ \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^m} - \frac{\partial \hat{g}_{lm}^{j*}}{\partial \hat{z}^k} \right\} B_l^i B_h^j B_l^k \\ + \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^k} \{ B_l^i B_h^k - B_h^i B_l^k \} \hat{z}_5^k B_l^m.$$

Transvecting (4.16) by \hat{z}_5^m and using (4.1) and (3.13) once more we find,

$$(4.17) \quad \left\{ \frac{\partial \hat{g}_{kl}^{j*}}{\partial \hat{z}^l} - \frac{\partial \hat{g}_{lh}^{j*}}{\partial \hat{z}^h} \right\} \hat{z}_5^m = \left\{ \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^m} - \frac{\partial \hat{g}_{lm}^{j*}}{\partial \hat{z}^k} \right\} \hat{z}_5^m B_l^i B_h^j \\ + g_{kl}^{j*} \{ B_l^i B_h^k - B_h^i B_l^k \} \hat{z}_5^k$$

while further multiplication by \hat{z}_5^l gives in view of (3.15),

$$(4.18) \quad \left\{ \frac{\partial \hat{g}_{kl}^{j*}}{\partial \hat{z}^l} - \frac{\partial \hat{g}_{lh}^{j*}}{\partial \hat{z}^h} \right\} \hat{z}_5^m \hat{z}_5^l = B_h^k \left\{ \frac{\partial \hat{g}_{km}^{j*}}{\partial \hat{z}^m} - \frac{\partial \hat{g}_{lm}^{j*}}{\partial \hat{z}^k} \right\} \hat{z}_5^m \hat{z}_5^l \\ + g_{kl}^{j*} B_l^i \hat{z}_5^k \hat{z}_5^l.$$

The relations (4.11) and (4.18) provide some indication as to how the covariant derivatives can be formed. In connection of first of these we have from (4.7),

$$(4.19) \quad g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m B_n^k = \left\{ g_{km}^{\alpha\sigma} B_n^k z_{\alpha\gamma}^m + g_{km}^{\alpha\sigma} B_{\gamma\alpha}^m z_n^k \right\} B_n^k.$$

To eliminate $B_{\gamma\alpha}^m$ from (4.19) and (4.11) we subtract the latter from former. Thus we get with the use of (4.8),

$$(4.20) \quad B_n^k \left\{ g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m + \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^{\alpha}} z_{\alpha\gamma}^m \right\} \\ = g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m + \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^{\alpha}} z_{\alpha\gamma}^m z_n^k.$$

This relation shows that the quantities defined by

$$(4.21) \quad Z_{\alpha\gamma}^k = g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m + \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^{\alpha}} z_{\alpha\gamma}^m z_n^k$$

are the components of a covariant vector.

Similarly it can be verified that the quantities defined by

$$(4.22) \quad Z_{\alpha\gamma}^{\alpha\sigma} = g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m + \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^k} z_{\alpha\gamma}^m z_n^k$$

satisfies the transformation law

$$(4.23) \quad Z_{\alpha\gamma}^{\alpha\sigma} = B_{\alpha}^{\alpha'} B_{\gamma}^{\gamma'} Z_{\alpha'\gamma'}^{\alpha\sigma}$$

so that (4.21) is to be taken in conjugation with (4.22).

These are the required relations resulting from (4.11).

On the other hand from (4.6) and (4.8) we have

$$(4.24) \quad g_{km}^{\alpha\sigma} B_{\gamma}^j z_{\alpha}^j B_n^k = g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m B_n^k - g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m.$$

This is substituted in the right hand side of (4.18).

In this manner we find that

$$(4.25) \quad \bar{X}_n = B_n^k X_k$$

where

$$(4.26) \quad X_k = g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m + \left\{ \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^k} - \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^{\alpha}} \right\} z_{\alpha\gamma}^m z_n^k.$$

It follows that the quantities thus defined also form the components of a covariant vector.

Similarly we can easily verify that the quantities defined by

$$(4.27) \quad X_{\alpha\sigma} = g_{km}^{\alpha\sigma} z_{\alpha\gamma}^m + \left\{ \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^{\alpha}} - \frac{\partial g_{km}^{\alpha\sigma}}{\partial z^k} \right\} z_{\alpha\gamma}^m z_n^k$$

satisfy the transformation law

$$(4.28) \quad \bar{X}_{i^*} = B_{i^*}^{k^*} X_{k^*}.$$

In the above we have found a way of constructing covariant vectors in terms of derivatives of \bar{z}_i^1 and $\bar{z}_i^{n^*}$ which indicates that there exist certain connection coefficients which are in the construction of these vectors. In order to obtain the explicit form of these coefficients we suppose that there exist quantities $g_{i^*}^{j^*}$ such that

$$(4.29) \quad g_{i^*}^{j^*} g_{j^*}^{k^*} = \delta_{i^*}^{k^*} \delta_{\gamma^*}^{\beta^*}.$$

The existence of (4.29) follows from the condition (D) of section 2. From (4.29) it also follows that

$$(4.30) \quad g_{i^*}^{j^*} g_{j^*}^{i^*} = \delta_{i^*}^{i^*} \delta_{\gamma^*}^{\beta^*}.$$

Also if in analogy to (4.2) we write

$$(4.31) \quad A_i^k = \frac{\partial \bar{z}^k}{\partial z^i}, \quad A_{i^*}^{k^*} = \frac{\partial \bar{z}^{k^*}}{\partial z^{i^*}}$$

and thereby obtain

$$(4.32) \quad g_{i^*}^{j^*} A_{j^*}^{k^*} = g_{i^*}^{j^*} B_{j^*}^{k^*}.$$

Multiplying (4.13) by \bar{z}_i^k , noting (4.1 and (3.9) we get

$$(4.33) \quad \frac{\partial \bar{g}_{i^*}^{j^*}}{\partial \bar{z}^i} \bar{z}_j^k = \frac{\partial g_{i^*}^{j^*}}{\partial z^i} \bar{z}_j^k B_{j^*}^{m^*} B_{i^*}^l + g_{i^*}^{m^*} B_{i^*}^k B_{j^*}^{m^*} \bar{z}_j^l.$$

Transvecting this relation by $g_{i^*}^{j^*} A_{j^*}^{k^*}$ and using (4.32) we get after some simplification

$$(4.34) \quad B_{i^*}^k \bar{z}_j^l = g_{i^*}^{j^*} \frac{\partial g_{i^*}^{k^*}}{\partial \bar{z}^i} \bar{z}_j^l B_{j^*}^{m^*} - g_{i^*}^{m^*} \frac{\partial g_{i^*}^{k^*}}{\partial z^p} \bar{z}_j^l B_{j^*}^p.$$

This relation shows that

$$(4.35) \quad B_{i^*}^k \bar{z}_j^l = B_{i^*}^l \bar{G}_{i^*, \gamma^*}^k - G_{p, \gamma^*}^k B_{j^*}^p$$

where

$$(4.36) \quad G_{p, \gamma^*}^k = g_{i^*}^{m^*} \frac{\partial g_{i^*}^{k^*}}{\partial z^p} \bar{z}_j^l B_{j^*}^l$$

and a corresponding relation on $\bar{G}_{i^*, \gamma^*}^k$ in the barred system. From (4.36) and (4.30) we can deduce that

$$(4.37) \quad g_{i^*}^{m^*} G_{p, \gamma^*}^k = \frac{\partial g_{i^*}^{k^*}}{\partial z^p} \bar{z}_j^l B_{j^*}^l$$

so that the covariant vector (4.22) can be expressed in the form

$$(4.38) \quad Z_{k^* \gamma}^{\alpha^*} = g_{k^* \alpha^*}^{\beta^*} (\dot{z}_{\beta^* \gamma}^{\alpha^*} + G_{\beta^* \gamma}^{\alpha^*} z^{\beta^*}).$$

Now let us put

$$(4.39) \quad G_{kh, \gamma}^{\alpha^*} = \frac{\partial G_{k, \gamma}^{\alpha^*}}{\partial z^h}.$$

From (4.1) we can deduce that

$$(4.40) \quad \frac{\partial \dot{z}_{\beta^* \gamma}^{\alpha^*}}{\partial z^h} = \delta_h^{\alpha^*} \delta_{\beta^* \gamma}^{\alpha^*}.$$

Differentiating (4.35) with respect to $\dot{z}_{\beta^* \gamma}^{\alpha^*}$, using (4.39) and (4.40) we get

$$(4.41) \quad B_{\beta^* \gamma}^{\alpha^*} \delta_{\beta^* \gamma}^{\alpha^*} = B_{\beta^* \gamma}^{\alpha^*} \bar{G}_{\beta^* \gamma}^{\alpha^*} - G_{\beta^* \gamma}^{\alpha^*} B_{\beta^* \gamma}^{\alpha^*}.$$

In particular if we put $\beta = \gamma$ in (4.41) we get

$$(4.42) \quad B_{\beta^* \beta^*}^{\alpha^*} = B_{\beta^* \beta^*}^{\alpha^*} \bar{\Gamma}_{\beta^* \beta^*}^{\alpha^*} - \Gamma_{\beta^* \beta^*}^{\alpha^*} B_{\beta^* \beta^*}^{\alpha^*}$$

where

$$(4.43) \quad \bar{\Gamma}_{\beta^* \beta^*}^{\alpha^*} = \frac{1}{m} G_{\beta^* \beta^*}^{\alpha^*}.$$

The relation (4.42) shows that $\Gamma_{\beta^* \beta^*}^{\alpha^*}$ obey the transformation law of the connection coefficients. Accordingly we shall regard the quantities (4.43) as the connection coefficients of our manifold C_n .

Similarly we can show that the quantities defined by

$$(4.44) \quad \Gamma_{\beta^* \alpha^*}^{\alpha^*} = \frac{1}{m} G_{\beta^* \alpha^*}^{\alpha^*}$$

where

$$(4.45) \quad G_{\beta^* \alpha^*}^{\alpha^*} = \frac{\partial G_{\beta^* \alpha^*}^{\alpha^*}}{\partial z^{\alpha^*}}$$

and

$$(4.46) \quad G_{\beta^* \gamma}^{\alpha^*} = g_{\beta^* \alpha^*}^{\gamma^*} \frac{\partial g_{\gamma^* \alpha^*}^{\beta^*}}{\partial z^{\alpha^*}} z^{\beta^*}$$

satisfy the transformation law

$$(4.47) \quad B_{\beta^* \alpha^*}^{\alpha^*} = B_{\beta^* \alpha^*}^{\alpha^*} \bar{\Gamma}_{\beta^* \alpha^*}^{\alpha^*} - \Gamma_{\beta^* \alpha^*}^{\alpha^*} B_{\beta^* \alpha^*}^{\alpha^*}.$$

From (4.36), (4.39) and (4.43) it follows that

$$(4.48) \quad m \Gamma_{\beta^* \beta^*}^{\alpha^*} = \frac{\partial g_{\beta^* \alpha^*}^{\alpha^*}}{\partial z^{\beta^*}} \frac{\partial z^{\alpha^*}}{\partial z^{\beta^*}} z^{\beta^*} + \frac{g_{\beta^* \alpha^*}^{\alpha^*}}{g_{\beta^* \alpha^*}^{\alpha^*}} \frac{\partial g_{\beta^* \alpha^*}^{\alpha^*}}{\partial z^{\beta^*}}$$

which shows that the connection coefficient $\Gamma_{\rho\sigma}^{\mu}$ is explicitly derivable from the metric tensor and its various derivatives. Similarly $\Gamma_{\rho\sigma\alpha}^{\mu}$ is derivable from the metric tensor and its derivatives.

5. - PARTIAL COVARIANT DERIVATIVES

We shall now use the connection coefficients defined by (4.43) and (4.44) to construct partial covariant derivatives of a given vector field $V_{\underline{c}}^{\underline{a}}$ ($x^i, x^{\alpha}, \dot{z}_{\lambda}^i, \dot{z}_{\lambda}^{\alpha}$) which is such that

$$(5.1) \quad V_{\underline{c}}^{\underline{a}}(x^i, x^{\alpha}, \dot{z}_{\lambda}^i, \dot{z}_{\lambda}^{\alpha}) = B_{\underline{c}}^{\underline{a}} \bar{V}_{\underline{c}}^{\underline{a}}(\bar{x}^i, \bar{x}^{\alpha}, \dot{\bar{z}}_{\lambda}^i, \dot{\bar{z}}_{\lambda}^{\alpha}).$$

Differentiation of this with respect to \dot{z}_{λ}^i , and $\dot{z}_{\lambda}^{\alpha}$ respectively gives

$$(5.2) \quad \frac{\partial V_{\underline{c}}^{\underline{a}}}{\partial \dot{z}_{\lambda}^i} = B_{\underline{c}}^{\underline{a}} \frac{\partial \bar{V}_{\underline{c}}^{\underline{a}}}{\partial \dot{\bar{z}}_{\lambda}^i} \Lambda_i^{\underline{a}},$$

$$(5.3) \quad \frac{\partial V_{\underline{c}}^{\underline{a}}}{\partial \dot{z}_{\lambda}^{\alpha}} = B_{\underline{c}}^{\underline{a}} \frac{\partial \bar{V}_{\underline{c}}^{\underline{a}}}{\partial \dot{\bar{z}}_{\lambda}^{\alpha}} \Lambda_{\alpha}^{\underline{a}}.$$

From (4.1) and (4.31) it also follow that

$$(5.4) \quad \frac{\partial \dot{z}_{\lambda}^i}{\partial x^j} = A_{kj}^i \dot{z}_{\lambda}^k, \quad \frac{\partial \dot{z}_{\lambda}^{\alpha}}{\partial x^{\beta}} = A_{k\beta}^{\alpha} \dot{z}_{\lambda}^k$$

where

$$(5.5) \quad A_{kj}^i = \frac{\partial^2 \bar{x}^i}{\partial x^k \partial x^j}, \quad A_{k\beta}^{\alpha} = \frac{\partial^2 \bar{x}^{\alpha}}{\partial x^k \partial x^{\beta}}$$

and since (4.31) is the inverse of (4.2) ie.

$$(5.6) \quad B_{\underline{a}}^{\underline{b}} \Lambda_{\underline{b}}^{\underline{a}} = \delta_{\underline{a}}^{\underline{a}}, \quad B_{\alpha}^{\beta} \Lambda_{\beta}^{\alpha} = \delta_{\alpha}^{\alpha}$$

we have

$$(5.7) \quad \Lambda_{kj}^i = -A_{in}^i B_{\rho\lambda}^n \Lambda_{\rho}^i \Lambda_{\lambda}^j; \quad \Lambda_{k\beta}^{\alpha} = -A_{\mu\sigma}^{\alpha} B_{\rho\lambda}^{\mu} \Lambda_{\rho}^{\alpha} \Lambda_{\lambda}^{\beta}.$$

Substituting (5.7) in (5.4) it follows that

$$(5.8) \quad \frac{\partial \dot{z}_{\lambda}^i}{\partial x^j} = -B_{\rho\lambda}^n \dot{z}_{\lambda}^{\rho} \Lambda_{\rho}^i \Lambda_{\lambda}^j; \quad \frac{\partial \dot{z}_{\lambda}^{\alpha}}{\partial x^{\beta}} = -B_{\rho\lambda}^{\mu} \dot{z}_{\lambda}^{\rho} \Lambda_{\rho}^{\alpha} \Lambda_{\lambda}^{\beta}.$$

Now differentiating (5.1) with respect to x^k , we obtain

$$(5.9) \quad \frac{\partial V_{\underline{c}}^{\underline{a}}}{\partial x^k} = B_{\underline{c}}^{\underline{a}} \Lambda_k^{\underline{a}} \bar{V}_{\underline{c}}^{\underline{a}} + B_{\underline{c}}^{\underline{a}} \left\{ \frac{\partial \bar{V}_{\underline{c}}^{\underline{a}}}{\partial \bar{x}^i} \Lambda_i^{\underline{a}} + \frac{\partial \bar{V}_{\underline{c}}^{\underline{a}}}{\partial \bar{x}^{\alpha}} \frac{\partial \dot{\bar{z}}_{\lambda}^{\alpha}}{\partial x^k} \right\}.$$

In the last term of the right hand side of this equation we substitute the value of $\delta \dot{z}_i^j / \delta z^k$ from (5.8) and then B_{pi}^m are eliminated with the help of (4.42). This process gives the following expression

$$- B_h^l \frac{\delta \bar{V}_e^h}{\delta \dot{z}_k^l} A_m^l A_k^m (B_i^m \bar{\Gamma}_{pi}^m - \Gamma_{pi}^m B_p^l B_q^l) \dot{z}_i^q$$

and by means of (5.2) we can easily reduce it to

$$- B_h^l A_k^q \frac{\delta \bar{V}_e^h}{\delta \dot{z}_k^l} \bar{\Gamma}_{pi}^q \dot{z}_i^p + \frac{\delta V_e^l}{\delta \dot{z}_k^l} \Gamma_{ik}^l \dot{z}_i^k.$$

Thus (5.9) can be written in the form

$$(5.10) \quad \frac{\delta V_e^l}{\delta z^k} - \frac{\delta V_e^l}{\delta \dot{z}_k^l} \Gamma_{ik}^l \dot{z}_i^k = B_h^l A_k^q \bar{V}_e^h \\ + B_h^l A_k^q \left\{ \frac{\delta \bar{V}_e^h}{\delta z^q} - \frac{\delta \bar{V}_e^h}{\delta \dot{z}_k^l} \bar{\Gamma}_{pi}^q \dot{z}_i^p \right\}.$$

By means of (4.42) the first term on the right hand side of this relation can be expressed as

$$A_k^l B_i^l \bar{\Gamma}_{li}^l \bar{V}_e^l - \Gamma_{kp}^l V_e^p.$$

Substituting this in (5.10) we get,

$$(5.11) \quad V_{e|k}^l = B_h^l A_k^q \bar{V}_{e|q}^h,$$

where

$$(5.12) \quad \bar{V}_{e|k}^h = \frac{\delta V_e^h}{\delta z^k} - \frac{\delta V_e^h}{\delta \dot{z}_k^l} \Gamma_{lk}^m \dot{z}_l^m + \Gamma_{lk}^h V_e^l.$$

The relation (5.11) shows that the quantities $V_{e|k}^l$ defined by (5.12) represent components of a mixed tensor. Thus $V_{e|k}^l$ represents a covariant partial derivative of the quantities $V_e^l (z^i, z^p, \dot{z}_i^j, \dot{z}_{i^*}^{j^*})$ with respect to z^k .

Now we shall find the covariant partial derivative of V_e^l with respect to z^{k^*} . For this purpose let us differentiate (5.1) with respect to z^{k^*} . Thus

$$(5.13) \quad \frac{\delta V_e^l}{\delta z^{k^*}} = B_h^l \left\{ \frac{\delta \bar{V}_e^h}{\delta z^{i^*}} A_{k^*}^{i^*} + \frac{\delta \bar{V}_e^h}{\delta \dot{z}_{i^*}^j} \frac{\delta \dot{z}_{i^*}^j}{\delta z^{k^*}} \right\}$$

Substituting the value of $\delta \dot{z}_{i^*}^j / \delta z^{k^*}$ from (5.8) and then eliminating $B_{p^*q^*}^m$ with the help of (4.47) we get (after using (5.3))

$$(5.14) \quad \frac{\delta V_e^l}{\delta z^{k^*}} - \frac{\delta V_e^l}{\delta \dot{z}_{i^*}^j} \Gamma_{i^*k^*}^m \dot{z}_{i^*}^{j^*} \\ = B_h^l A_{k^*}^{i^*} \left\{ \frac{\delta \bar{V}_e^h}{\delta z^{i^*}} - \frac{\delta \bar{V}_e^h}{\delta \dot{z}_{i^*}^j} \bar{\Gamma}_{p^*q^*}^m \dot{z}_{i^*}^{j^*} \right\}.$$

This relation shows that the quantities defined by

$$(5.15) \quad V_{\bar{c}|k}^{\bar{c}} = \frac{\partial V_{\bar{c}}^{\bar{c}}}{\partial z^{k*}} - \frac{\partial V_{\bar{c}}^{\bar{c}}}{\partial \bar{z}_{\lambda}^{\bar{c}}} \Gamma_{p k}^{* \bar{c}} \bar{z}_{\lambda}^{* \bar{c}}$$

represent a component of a mixed tensor. Thus $V_{\bar{c}|k}^{\bar{c}}$ represent a covariant partial derivative of $V_{\bar{c}}^{\bar{c}}$ with respect to z^{k*} .

In the similar manner we can construct the covariant partial derivative of a vector field $V_{\bar{c}}^{\bar{c}}$ ($z^{\bar{c}}, z^{k*}, \bar{z}_{\lambda}^{\bar{c}}, \bar{z}_{\lambda}^{* \bar{c}}$) with respect to $z^{\bar{c}}$ and z^{k*} respectively. These are given by

$$(5.16) \quad V_{\bar{c}|k}^{\bar{c}} = \frac{\partial V_{\bar{c}}^{\bar{c}}}{\partial z^{\bar{c}}} - \frac{\partial V_{\bar{c}}^{\bar{c}}}{\partial \bar{z}_{\lambda}^{\bar{c}}} \Gamma_{p k}^{\bar{c}} \bar{z}_{\lambda}^{\bar{c}},$$

$$(5.17) \quad V_{\bar{c}|k}^{* \bar{c}} = \frac{\partial V_{\bar{c}}^{* \bar{c}}}{\partial z^{k*}} - \frac{\partial V_{\bar{c}}^{* \bar{c}}}{\partial \bar{z}_{\lambda}^{* \bar{c}}} \Gamma_{p k}^{* \bar{c}} \bar{z}_{\lambda}^{* \bar{c}} + \Gamma_{k^{* \bar{c}}}^{\bar{c}} V_{\bar{c}}^{* \bar{c}}.$$

The existence of connection coefficients and covariant partial derivatives of a vector field suggests the existence of curvature tensors and a corresponding theory of curvature. We hope to deal the brief discussion of these concepts in a separate publication.

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